

Emergent Electroweak Symmetry Breaking

Yanou Cui

Harvard University

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with Tony Gherghetta, James Stokes (university of Melbourne)
and James D. Wells (University of Michigan & CERN)

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Outline

- 1 Motivation
- 2 Model Setup (W,Z, γ)
- 3 Electroweak Precision Test
- 4 Fermion Sector
- 5 Phenomenology
- 6 Conclusions

Motivation

- **Standard Model:** Benchmark model of particle physics, based on gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y$
 Story incomplete: W, Z bosons, fermions are **MASSIVE** →
 EW gauge symmetry not exact
Big mystery to unveil at the LHC: **Origin of masses for SM particles?**
- Most existing solutions: mass origin *equivalent to* **(fundamental) Electroweak gauge symmetry breaking (EWSB)** via Higgs mechanism
- **Higgs Mechanism:** ♣; Higgs boson: in hot pursuit (but not seen yet...)
Is there alternative clue to 'mass origin' beyond this paradigm? Surprises from Nature?

A brief Review of existing EWSB models

- **Elementary Higgs:** EW scale stabilized by SUSY, EWSB triggered by dynamical SUSY breaking in a strong hidden sector
- **Composite Higgs (5D dual):** pseudo-Goldstone boson of chiral symmetry breaking
- **Technicolor-like Higgsless (5D dual):** decoupled heavy σ mode, strong TC sector

Common features of these EWSB models (4D)

- To naturally solve ‘gauge hierarchy’ problem—i.e. **generate** a TeV mass gap via dimensional transmutation, require a new **external** sector **beyond** the SM with **confining strong dynamics**
- Most of **SM fields**, esp. gauge fields stay **elementary**, **spectators** of strong dynamics, acquire mass by coupling to the strong sector via **Higgs mechanism**

- A ‘wild’ curiosity: Why this strong dynamics almost ‘inevitable’ for generating SM mass has to be **external**? Why not **underlying** the SM?
- **Actually not ‘wild’**: familiar example—QCD! Mesons, baryons—**composites** getting mass *directly* from quark **confinement**; esp. **vector ρ meson** acts like a massive ‘**gauge field**’ of spontaneously broken **hidden local symmetry** (more later...)
- **Different origins of mass**:
 - Higgs mechanism: W, Z slowed down when ‘swimming’ through dark medium-Higgs (*‘cosmic superconductor’*)
 - Compositeness: W, Z gain mass from binding energy of internal constituents, like tension of a string

⇒ An **alternative** origin of mass where ‘**EWSB**’ is IR Emergent, opens up **deeper substructure in Nature**?

More on emergent gauge theory

- Emergent gauge symmetry (breaking) (or composite gauge boson) is **conceptually innocuous, even inspiring**:
'Gauge symmetry, unlike global symmetry, is not a true symmetry of nature, does not lead to new conserved charge, merely reflect redundancy in the description' (D. J. Gross, "Gauge theory - past, present and future," 1997)
- **Theoretical example**—Seiberg duality: in conformal window $\mathcal{N} = 1$ SUSY QCD, a dual 'magnetic gauge theory' $SU(N_f - N_c)$ *emerges at composite level*, with same IR physics as original $SU(N_c)$ 'electric gauge theory'
"gauge symmetry may not be fundamental, some gauge symmetries in the SM or even general relativity may be long distance artifacts." (N. Seiberg, "The power of duality: Exact results in 4D SUSY field theory,"1995)
- **Realistic example**—Hidden local symmetry (HLS) in QCD→

Intriguing story of QCD ρ meson

QCD ρ meson: vector boson in $SU(2)_V$ triplet, ρ^\pm, ρ^0 obtain almost degenerate masses from QCD confinement.

Phenomenology: Universal coupling to matter ($\pi, N\dots$); ρ -dominance of hadron EM form factor and of $\pi\pi$ scattering

A neat theory to describe all these observations: ρ meson is a dynamical gauge boson of spontaneously-broken HLS $SU(2)_V$ (Sakurai first conjectured in 1960, M. Bando etc. 1985)

Remarkable similarity to W, Z in EW theory

- Unbroken QCD $SU(2)_V$ (global) acts like **custodial symmetry** ensures degenerate $\rho^{1,2,3}$ masses with no $\rho - \gamma$ mixing
- $\rho - \gamma$ mixing is in exactly the same pattern as $W^3 - B$ mixing in EW theory ($SU(2)_{HL} \times U(1)_Q \rightarrow U(1)'$) and splits ρ^0, ρ^\pm masses

Could EWSB be another example of 'emergent gauge theory'? W,Z composites?

- **History:** Early attempts— Bjorken 1977; Hung, Sakurai 1978; Abott, Farhi 1981; Fritzsche, Schildnecht, Kogerlert 1982; Kugo, Uehara, Yanagida 1984; Suzuki 1987, [arXiv:1006.1319](#)...
Model Building, calculability limited due to nonperturbative nature
- **Our Goal:** towards a **realistic** composite W, Z model with **good calculability** using powerful modern tool— **AdS/CFT**: 5D physics in warped Xdim \Rightarrow Effective 4D chiral \mathcal{L} of massive W, Z, fermions
- **Job list:**
 - Generate masses for W, Z, fermions
 - Ensure universality of gauge couplings
 - EW precision tests
 - Novel solution for unitarizing high E $W_L W_L$ scattering: form factor suppression for composites...

Organizing Principles, AdS/CFT Dictionary

- AdS/CFT correspondence (Maldacena, 1997):
strongly-coupled 4D conformal gauge theory \Leftrightarrow weakly coupled 5D (AdS_5) gravity theory
- CFT with a UV cutoff, undergoes CFT-breaking confinement at IR \Leftrightarrow Warped 5D model (e.g. RS1) built on a slice of AdS_5 (Arkani-Hamed, Porrati, Randall 2001 etc.)
- Global symmetry (\rightarrow emergent gauge symmetry at IR) in 4D CFT, unbroken at IR \Leftrightarrow 5D gauge symmetry broken on UV brane, preserved on IR brane
Composite gauge boson \Leftrightarrow KK mode of 5D gauge field peaking on IR brane

Two points to keep in mind:

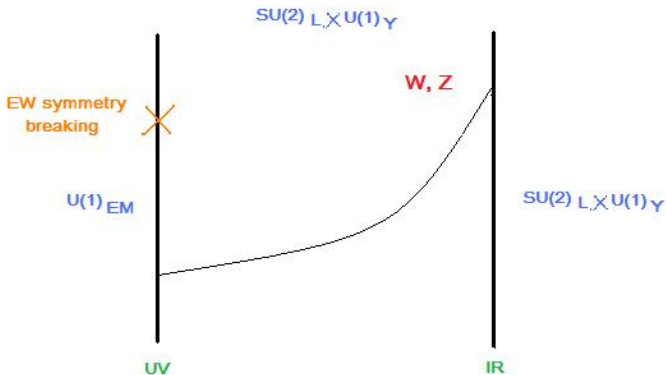
- To ensure a **robust** weakly coupled gravity dual, **implicitly assume** the 4D side is **large N, large t' Hooft coupling** (unlike QCD! \Rightarrow possible interesting phenomenology **later...**)
- **5D EW gauge symmetry broken at UV**: different from original RS and all existing warped models where EWSB occurs on IR
Simple picture: 'would-be' elementary gauge field (flat zero mode) get Planck scale mass upon EWSB at UV, decouples from low energy theory; TeV physics described by light KK modes, W, Z are lightest KK, mass do not come from 4D Higgs mechanism

Getting realistic mass spectrum

- Challenge:** composite W,Z \leftrightarrow 1st KK modes peaking at IR; Usual KK mass spectrum in RS1-type model $m_n \sim n z_1^{-1} \Rightarrow m_Z = m_1 \sim 90\text{GeV}$ implies $m_2 \sim 200\text{GeV} \leftarrow$ LEP bound $m_{Z'} > 1500\text{GeV}$
 \Rightarrow (On 4D side) W, Z needs to be 'special light mode' vs. confinement or typical resonance scale at O(TeV) . Known e.g. QCD $m_\pi \ll \Lambda_{QCD}$
- How to realize in 5D model—a 'distorted' spectrum with ultra-light 1st KK?—Turn on **brane kinetic term (BKT)**: Generic, compatible with 4D Poincare symmetry on brane
 — (Carena, Ponton, Tait and Wagner 2003; Davoudiasl, Hewett and Rizzo 2003)

The 5D Model

Model Setup:



5D action

$$\begin{aligned}
 S &= \int d^4x dz \sqrt{-g} \left[-\frac{1}{4} (F_{MN}^{La})^2 - \frac{1}{4} (F_{MN}^Y)^2 \right. \\
 &\quad - \frac{1}{2} (kz) \delta(z - z_{UV}) \frac{\zeta_Q}{g_{Y5}^2 + g_{L5}^2} (g_{Y5} F_{\mu\nu}^{L3} + g_{L5} F_{\mu\nu}^Y)^2 \\
 &\quad \left. - \frac{1}{2} (kz) \delta(z - z_{IR}) (\zeta_L (F_{\mu\nu}^{La})^2 + \zeta_Y (F_{\mu\nu}^Y)^2) \right].
 \end{aligned}$$

Boundary conditions

$$\begin{aligned}
 z = z_{UV} : & \begin{cases} \partial_z (g_{Y5} A_\mu^{L3} + g_{L5} B_\mu) + \zeta_Q \square (g_{Y5} A_\mu^{L3} + g_{L5} B_\mu) = 0, \\ g_{L5} A_\mu^{L3} - g_{Y5} B_\mu = 0, \\ A_\mu^{L1,2} = 0, \end{cases} \\
 z = z_{IR} : & \begin{cases} \partial_z A_\mu^{La} - \zeta_L k z_{IR} \square A_\mu^{La} = 0, \\ \partial_z B_\mu - \zeta_Y k z_{IR} \square B_\mu = 0, \end{cases}
 \end{aligned}$$

Mass spectrum and Profiles

- A flat zero mode exists: photon
- Lightest KK mode: SM W, Z boson

$$m_Z \simeq \sqrt{\frac{2}{\zeta_{Lk}} + \frac{2}{\zeta_{Qk}(1+\beta^2)}} z_{IR}^{-1}.$$

$$m_W \simeq \sqrt{\frac{2}{\zeta_{Lk}}} z_{IR}^{-1}.$$

- Higher KK modes have 'normal' masses of $O(z_{IR}^{-1})$ (TeV), suppressed IR coupling due to BKT induced repelling (Carena, Ponton, Tait and Wagner 2003; Davoudiasl, Hewett and Rizzo 2003)

W, Z 5D profiles:

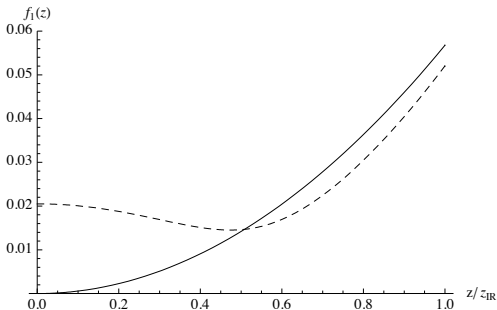


Figure: The W -boson (solid) and Z -boson (dashed) profiles in units of \sqrt{k} .—Peaking at IR ($\sim \propto z^2$), indeed dual composites

Electroweak Precision Test

T parameter: deviation from relation $\rho \equiv \frac{m_W^2}{m_Z^2 \cos^2 \theta_w} = 1$

- Sufficient protection requires **custodial symmetry $SU(2)$** $\Rightarrow m_W^\pm = m_Z$ at $g_1 = 0$, Automatic in SM Higgs with potential $SO(4)$
- Most new models need to add global $SU(2)_R$ **in addition to** $SU(2)_L \times U(1)_Y$ (e.g. additional bulk $SU(2)$ in most RS models)

Our model: No need for new $SU(2)$, $SU(2)_L$ itself play the role!

- **Our W,Z are true KK modes**, masses correlated by **bulk $SU(2)_L$** upon dim-reduction
- In 4D dual picture, masses originate from CFT breaking at IR which **breaks 'emergent' local $SU(2)$** , while BC ensures **global $SU(2)_L$ unbroken** there
- **Recall ρ meson in QCD!**—upon confinement $SU(2)_V$ global survives, emergent $SU(2)$ local is 'broken' since triplet ρ get masses (degenerate due to $SU(2)_V$ global)

S-parameter: new physics contribution to Π'_{33}, Π'_{3Q}

- **A challenge for Higgsless models:** $S \sim \frac{g_n}{M_n^2}$ (M_n : KK mass, g_n : KK V'-SM V coupling), in TC or its 5D variations typically g_n big, M_n relatively light $\Rightarrow S \sim O(1)$, too big
- **In our model:** BKT effects increases m_{KK} , decreases $g_n \Rightarrow$ reduce S

Matching to SM: Assumption for fermions: predominantly on IR brane (details on fermion sector, see later slides)

Calculation of S, T

- By definition $S \equiv 16\pi(\Pi'_{33} - \Pi'_{3Q})$, in 5D model gives

$$S = \frac{16\pi}{g^2 + g'^2} (1 - Z_Z) \text{ where}$$

$$Z_Z = \left\{ \int \frac{dz}{kz} \left[(f_1^{L3}(z))^2 + (f_1^B(z))^2 \right] + \frac{\zeta_Q}{1 + \beta_5^2} \left[f_1^{L3}(z_{UV}) + \beta_5 f_1^B(z_{UV}) \right]^2 \right. \\ \left. + \zeta_L (f_1^{L3}(z_{IR}))^2 + \zeta_Y (f_1^B(z_{IR}))^2 \right\},$$

- By definition $T \equiv \frac{4\pi}{\sin^2 \theta_w \cos^2 \theta_w m_Z^2} (\Pi_{11}(0) - \Pi_{33}(0))$, in 5D model

$$g^2 \Pi_{11}(0) = \int \frac{dz}{kz} |\partial_z f_W(z)|^2,$$

$$(g^2 + g'^2) \Pi_{33}(0) = \int \frac{dz}{kz} \left[|\partial_z f_1^{L3}(z)|^2 + |\partial_z f_1^B(z)|^2 \right].$$

Numerically find with $z_{IR} = 1.8\text{TeV}$,

$S \simeq 0.1$, $T \simeq 0.05$ —good fit with LEP EWPT data, confirm the built-in protection mechanisms discussed earlier

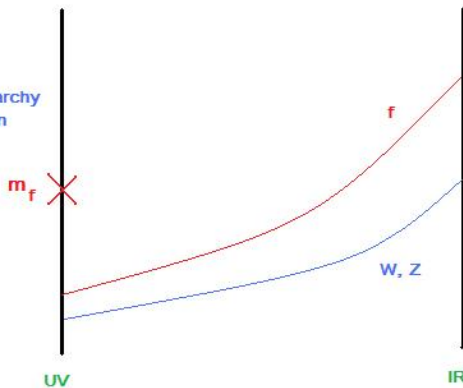
Fermion Sector

Guidelines:

- Gauge coupling universality:** usually automatic with flat gauge profile (elementary) once fermions canonically normalized
 Our model: **non-flat gauge profile** \Rightarrow need **universal fermion profile, same bulk mass c** , \Rightarrow fermion mass hierarchy cannot result from 'geography'
- Mass origin:** hierarchical UV Yukawa (Froggatt-Nielsen)+ **exponential suppression** of UV profile
 Fermion zero mode can have very different profile depending on bulk mass c : $\psi(y)^{(0)} \propto e^{(\frac{1}{2}-c)ky}$
 Naive estimation of mass from UV Yukawa y :
 $m \sim y \cdot k(\psi(y_{UV})^{(0)})^2 \sim y \cdot ke^{2(c-1/2)ky_{IR}} \Rightarrow$ with $c = 0$, $y \sim O(1)$
 gives $m \sim \text{TeV (top)}$ –**Magic from warping!**
- Predominantly IR localized:** add IR BKT

Model Setup

Fermion mass hierarchy
via Froggatt-Nielsen
mechanism



5D action

$$S_\Psi = \int d^5x \sqrt{-g} \left[\frac{1}{2} (\bar{\Psi}_i^{(L)} \Gamma^M D_M \Psi_i^{(L)} - D_M \bar{\Psi}_i^{(L)} \Gamma^M \Psi_i^{(L)}) + m_L^{(i)} \bar{\Psi}_i^{(L)} \Psi_i^{(L)} + (L \leftrightarrow R) \right].$$

$$S_m^{(UV)} = \int d^5x \sqrt{-g} \lambda_5^{(i)} \left[\bar{\Psi}_i^{(L)} \Psi_i^{(R)} + \bar{\Psi}_i^{(R)} \Psi_i^{(L)} \right] (kz) \delta(z - z_{UV}),$$

$$S_{KE}^{(IR)} = \int d^5x \sqrt{-g} \left[\frac{1}{2} \eta_{iL} (\bar{\Psi}_i^{(L)} \Gamma^\mu D_\mu \Psi_i^{(L)} - D_\mu \bar{\Psi}_i^{(L)} \Gamma^\mu \Psi_i^{(L)}) + (L \leftrightarrow R) \right] (kz) \delta(z - z_{IR})$$

Boundary Conditions

$$f_{L+}^{(n)}(z_{UV}^+) = -\lambda_5 f_{R+}^{(n)}(z_{UV}), \quad f_{L+}^{(n)}(z_{IR}^-) = -(\eta_{Lk})(m_n z_{IR}) f_{L-}^{(n)}(z_{IR}),$$

$$f_{R-}^{(n)}(z_{UV}^+) = \lambda_5 f_{L-}^{(n)}(z_{UV}), \quad f_{R-}^{(n)}(z_{IR}^-) = (\eta_{Rk})(m_n z_{IR}) f_{R+}^{(n)}(z_{IR}),$$

Example: Massless bulk fermions ($c = 0$)

Even modes profile:

$$|f_{L-}^{(n)}(z)| = |f_{R+}^{(n)}(z)| = N_n^{(0)} (kz)^2 [\cos(\hat{m}_n - m_n z) - (\eta k) \hat{m}_n \sin(\hat{m}_n - m_n z)]$$

where

$$N_n^{(0)} \simeq \frac{1}{\sqrt{z_{\text{IR}}}} \sqrt{\frac{1}{1 + (\eta k)/2 + (\eta k)^2 \hat{m}_n^2}}, \quad \hat{m} = m z_{\text{IR}}^{-1}$$

Obtain for BKT ($\eta = 10$):

$$m_e \leq m \leq m_t \quad \text{with} \quad 3.1 \times 10^{-6} \leq \lambda_5 \leq 1.15$$

W, Z Couplings

Fermion-gauge couplings determined by wavefunction overlap, e.g:

$$g_W \simeq \frac{g_{5L}}{\sqrt{\zeta_L}} \frac{2}{\eta k} \left[\frac{1}{3} + \xi_{IR} k (1 - 2(\eta k) \hat{m}_i \hat{m}_j) \right],$$

where BIT $\xi_{IR} = \sqrt{\zeta_L/2\Lambda\eta}$, $\hat{m} = mz_{IR}^{-1} \ll 1$

⇒ Light fermions: nonuniversality at the per-mille level

Top-related: nonuniversality at 15 – 25% level:

$$\frac{g_{W-(tb)}}{g_{W-(ud)}} = 0.854; \quad \frac{g_{Z-(top)}}{g_{Z-}^{(SM)}(top)} = 0.731; \quad \frac{g_{Z+(top)}}{g_{Z+}^{(SM)}(top)} = 0.732$$

Compatible with current direct bound:

Wtb: 20%@ Tevatron, Ztt: 40%@ LHC with 300fb^{-1}

Gauge coupling universality due to light fermion masses!

Top anomaly

Anomalous top couplings \Rightarrow **Indirect** bound from EWPT data

(Larios, Perez, Yuan '99)

Obtain:

$$\epsilon_1^{SM} + \delta\epsilon_1 \simeq 19 \times 10^{-3}; \quad \epsilon_b^{SM} + \delta\epsilon_b \simeq -13 \times 10^{-3},$$

This compares with (68% C.L.)

$$4.4 \times 10^{-3} \leq \epsilon_1^{exp} \leq 6.4 \times 10^{-3},$$
$$-6.2 \times 10^{-3} \leq \epsilon_b^{exp} \leq -3.1 \times 10^{-3}.$$

\Rightarrow Requires special treatment for top quark (common issue for Higgsless models)

e.g. separate brane for the top (Cacciapaglia, Csaki, Grojean, reece, Terning '05)

Some Interesting Phenomenology

- Anomalous top coupling; lightest KK- t' : 1.5TeV well in LHC reach ($2t+2\gamma/Z$ signal)
- **Compositeness** \Rightarrow **form factor** falling at high $p \Rightarrow$ Different story for WW scattering unitarization, anomalous WWZ , $WW\gamma$ couplings
- 4D dual theory: **large N**, **large t'Hooft coupling** \Rightarrow huge splitting rate \Rightarrow **No partons inside hadrons**, coherent scattering at DIS with p-falling form factor!
(Polchinski, Strassler 2002)

Conclusions

- We explore a new scenario where SM W, Z, fermions are **composites** getting masses from **strong dynamics** confining at IR, **without Higgs mechanism**. **SM EW gauge symmetry, EWSB are IR "emergent" phenomena**.
- Using AdS/CFT we build a calculable 5D warped model where EWSB on UV. \Rightarrow **realistic mass spectrum** & **good fit to S, T parameters**
- The composite nature of W, Z \Rightarrow **novel solution for WW scattering unitarization**; **coupling deviations from SM** can be probed at the LHC
- Indirect constraints: further model building required to successfully incorporate top

Novel Prospect at LHC: Mystery of SM mass origin \Rightarrow Deeper level of substructure in Nature?