

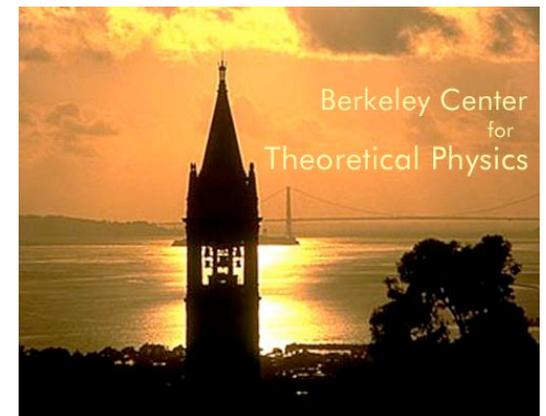


Jets and Jet Shapes in SCET

Christopher Lee

Lawrence Berkeley National Laboratory
and Berkeley Center for Theoretical Physics

Santa Fe 2010 Meeting
July 6, 2010



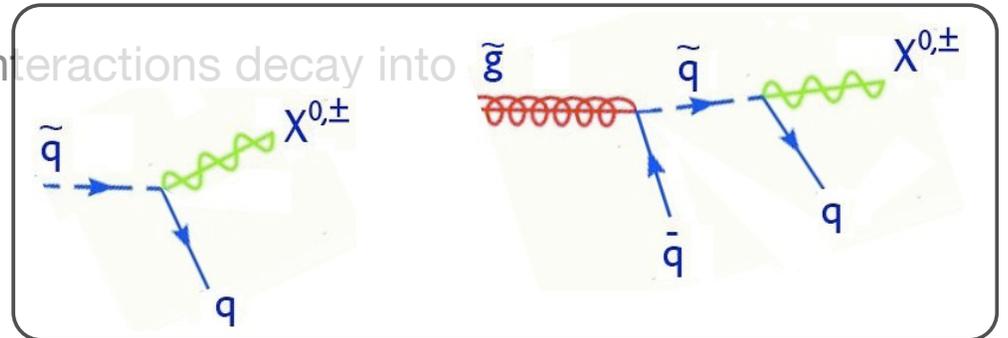
Jets as Signals of and Backgrounds to New Physics

- **Signal: New Physics** with strong interactions decay into **Jets**
 - e.g. Squark and Gluino Decays
- **Backgrounds:** Standard Model is very good at producing jets copiously!
- **Probing inside jets** yields more clues to the origin of jets, identifying underlying hard interaction
- **Important Standard Model applications**
 - test of understanding of perturbative vs. nonperturbative QCD
 - precision extractions of strong coupling constant from jet structure

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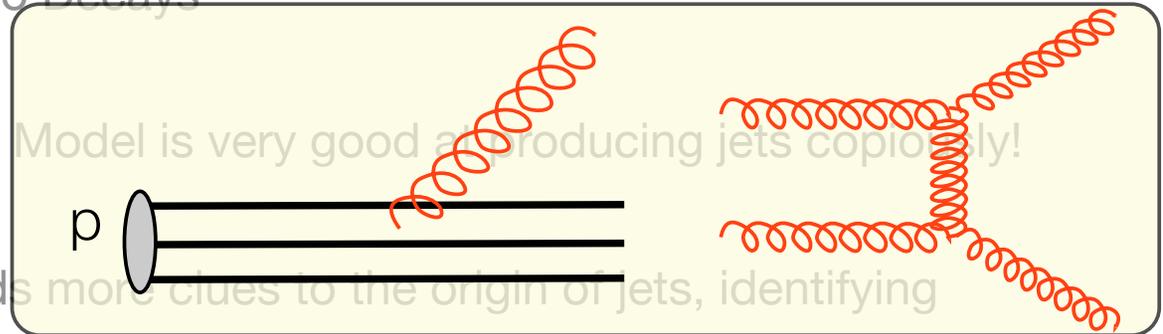
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Challenges in Predicting Jet Observables

- Example: Two-Jet Event Shapes

- Thrust:
$$T = \frac{1}{Q} \max_{\hat{t}} \sum_i |\hat{t} \cdot \mathbf{p}_i|$$

Farhi (1977)

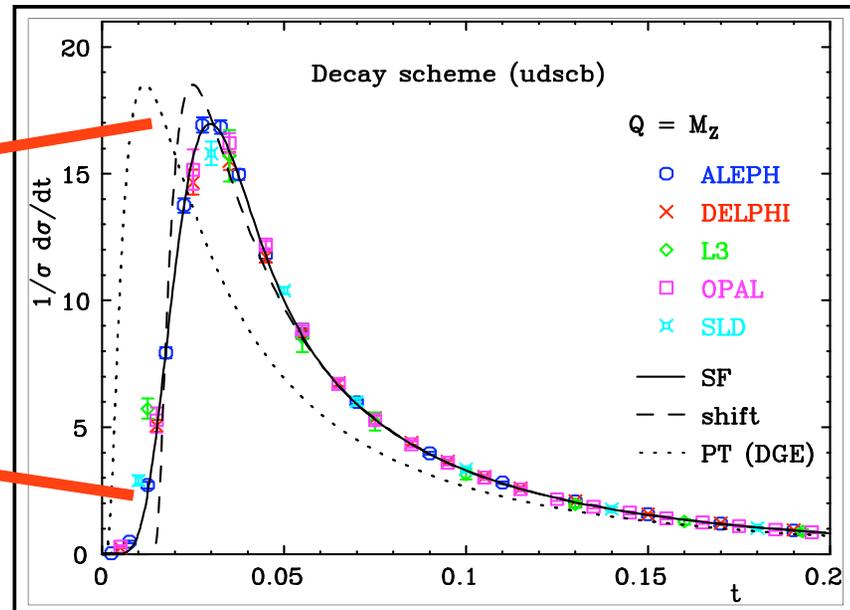
- Distribution of events in $\tau = 1-T$:



$\frac{1}{\tau Q}$ power corrections
due to nonperturbative
soft radiation

fixed-order PT blows
up, large logarithms
must be resummed to
obtain physical
behavior

$$\alpha_s^m \frac{\ln^n \tau}{\tau}$$



Gardi, Rathsmann (2002)

Theoretical Tools to Tackle Jet Physics

Tasks 	Tools 
Disentangling Perturbative and Nonperturbative Physics	<ul style="list-style-type: none">• Factorization<ul style="list-style-type: none">• Power counting• Separation of Scales
Resumming Large Logarithms in Perturbation Series	<ul style="list-style-type: none">• Renormalization Group Evolution<ul style="list-style-type: none">• QCD IR \longrightarrow EFT UV
Predictive Descriptions of Nonperturbative Effects	<ul style="list-style-type: none">• Universal Soft Matrix Elements<ul style="list-style-type: none">• e.g. PDFs, soft shape functions

Plan of the Talk

- A Theory Toolbox to Probe Jets: Soft-Collinear Effective Theory
 - Tools for Factorization: Power Counting and Soft-Collinear Decoupling
 - Tools for Resummation: the Renormalization Group
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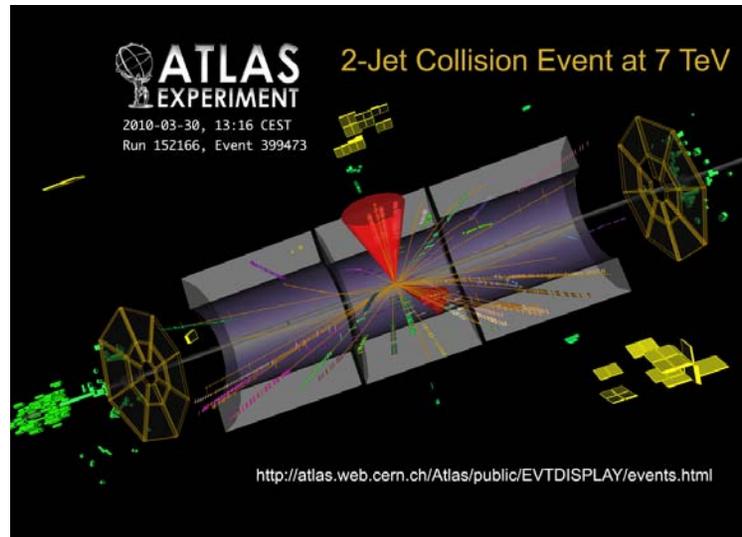
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Factorization of Jet Cross Sections

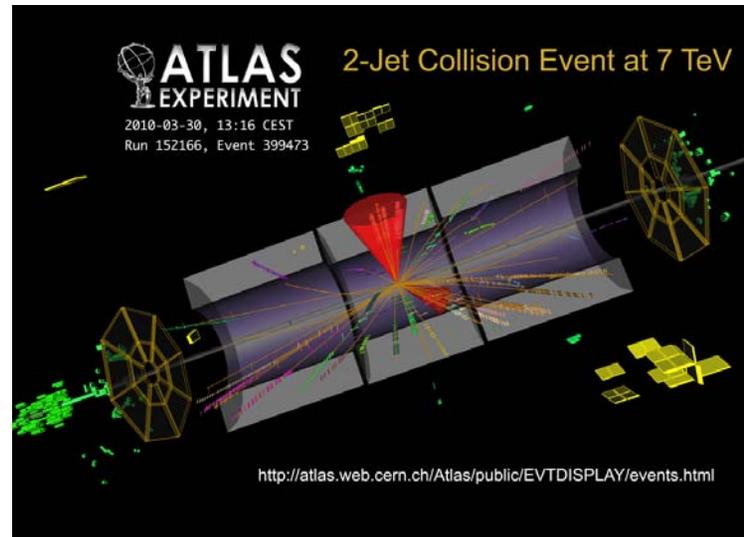
Collins, Soper, Sterman



Factorization of Jet Cross Sections

Collins, Soper, Sterman

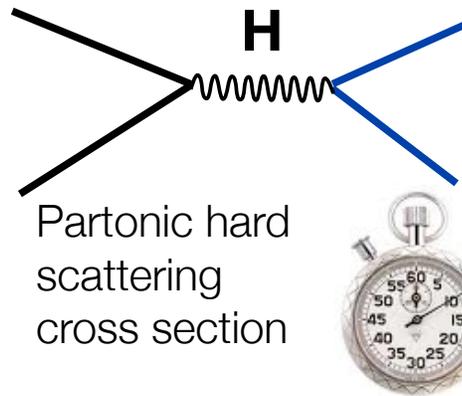
***Separate into physics at
disparate time scales***



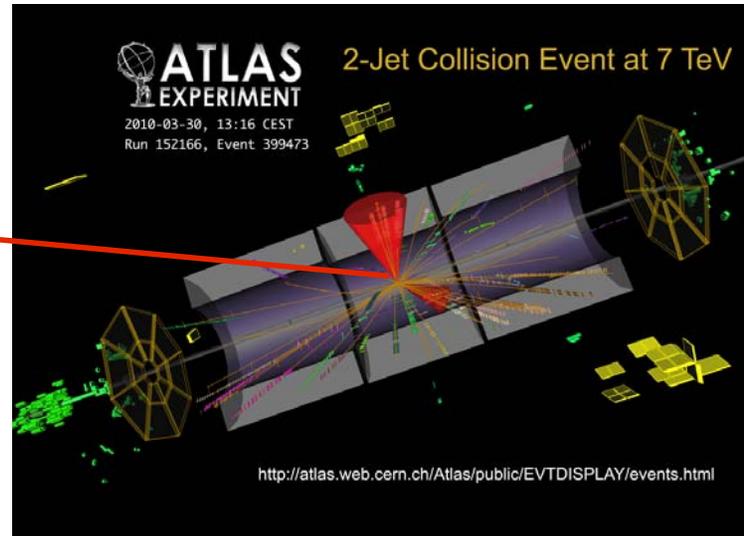
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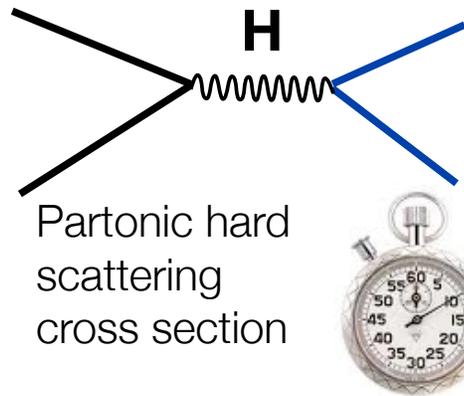
Partonic hard scattering cross section



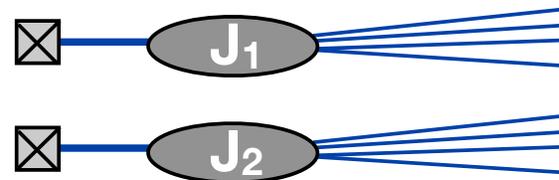
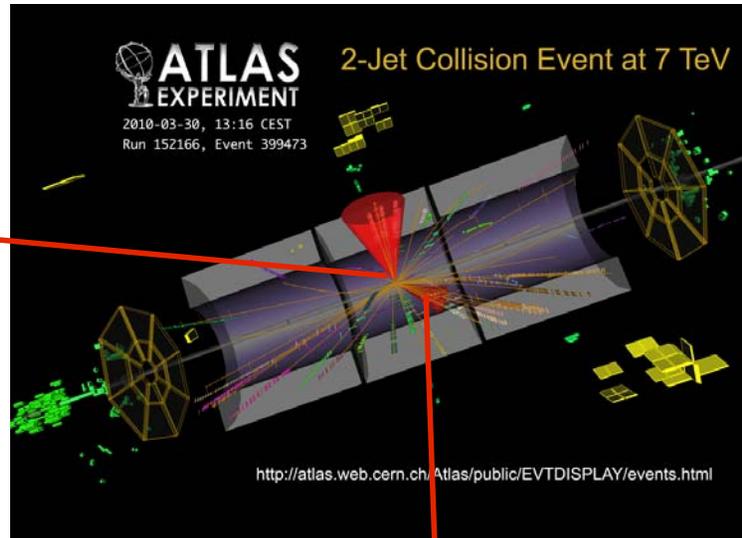
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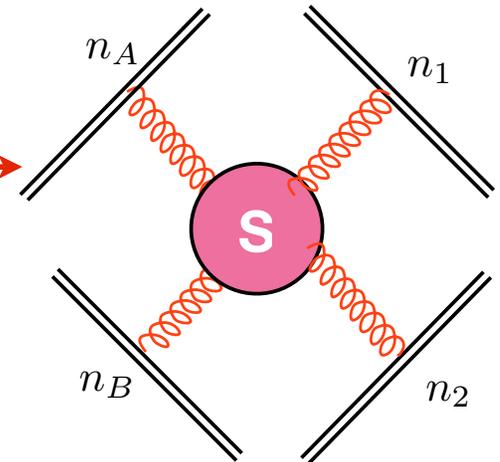
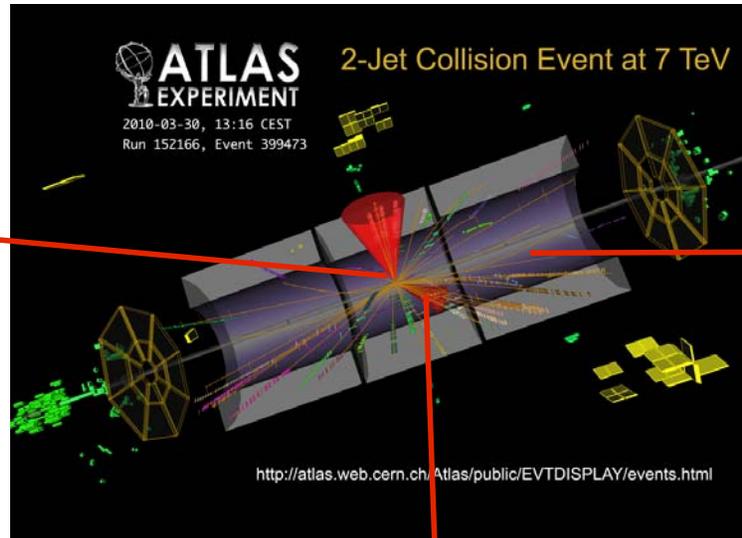
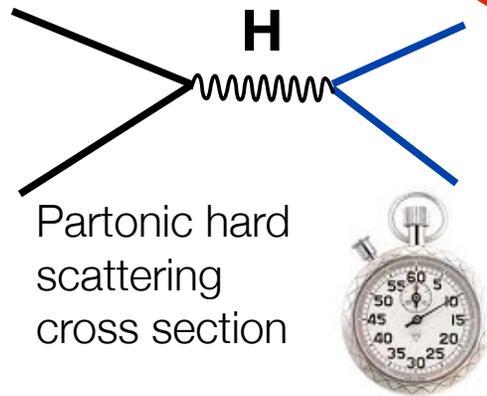
Collinear branching and showering
(*cf. parton shower*)



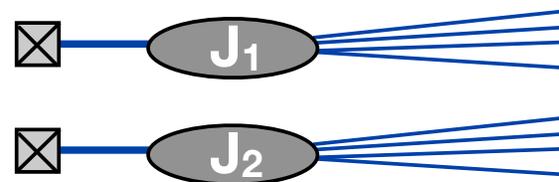
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Soft radiation, color exchange (hadronization)



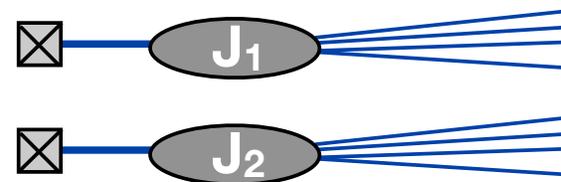
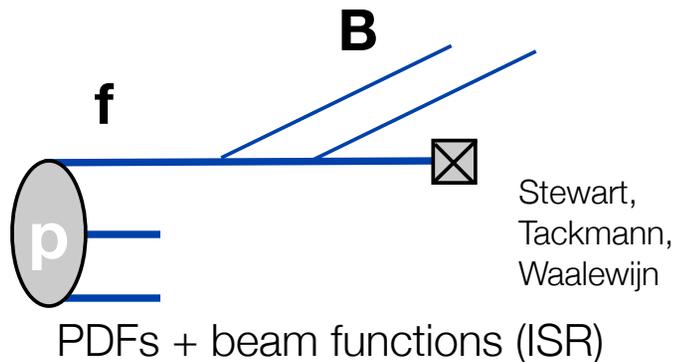
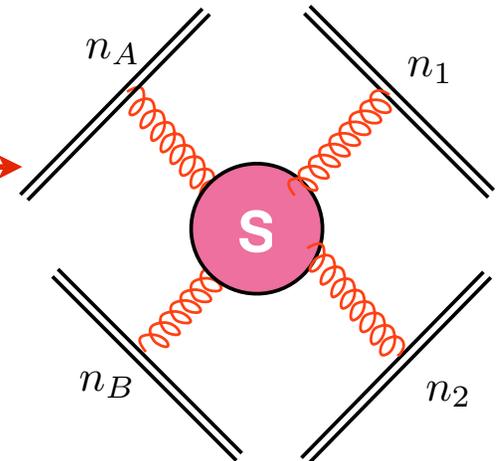
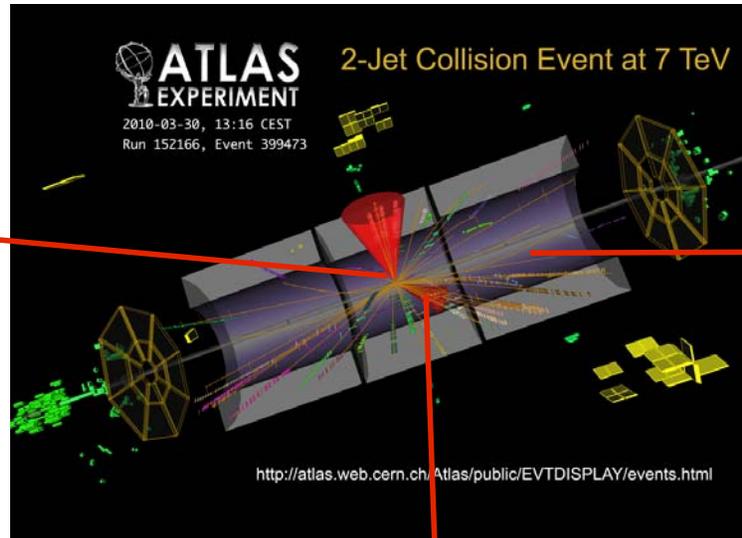
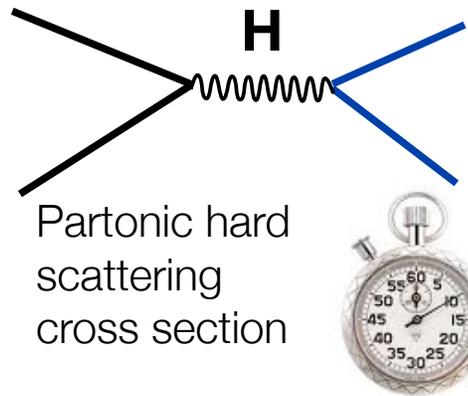
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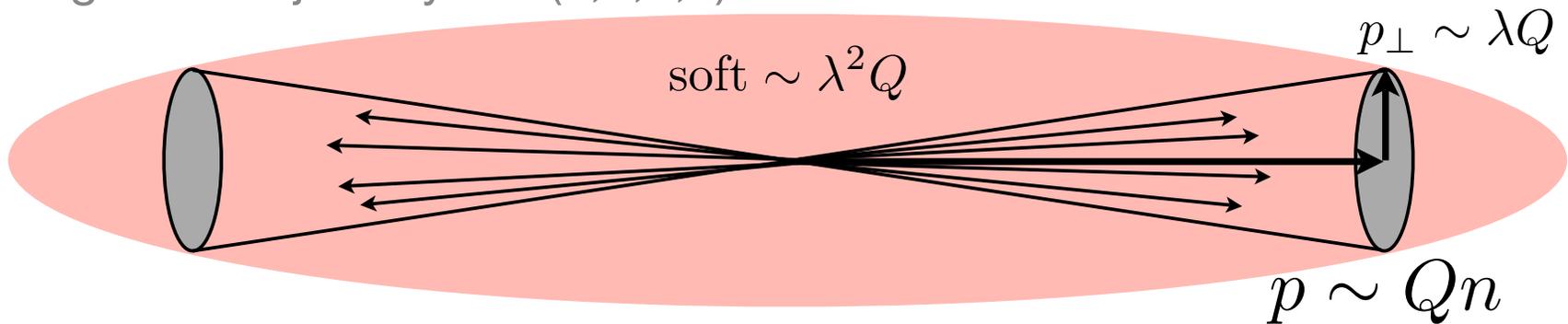
Effective Field Theory

- Use degrees of freedom relevant at the energy or distance scales at which one is working
- Expand full theory quantities in powers of small parameters, ratios of disparate scales
- **Effective theory often possesses enhanced symmetries or simplifications not immediately evident in full theory**
- **Renormalization group evolution between scales resums large logarithms in perturbative series**

Soft-Collinear Effective Theory

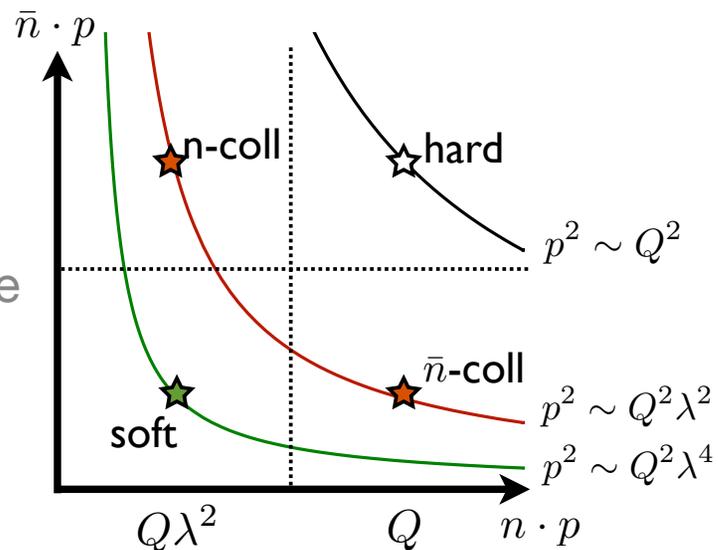
Bauer, Fleming, Luke,
Pirjol, Stewart (2000, 2001)

- Effective field theory for small momentum fluctuations about lightlike trajectory $n = (1,0,0,1)$



- In momentum space:

integrate out modes
of large virtuality one at a time



A Brief and Biased History of SCET

2000

2001

2003

2005

2006

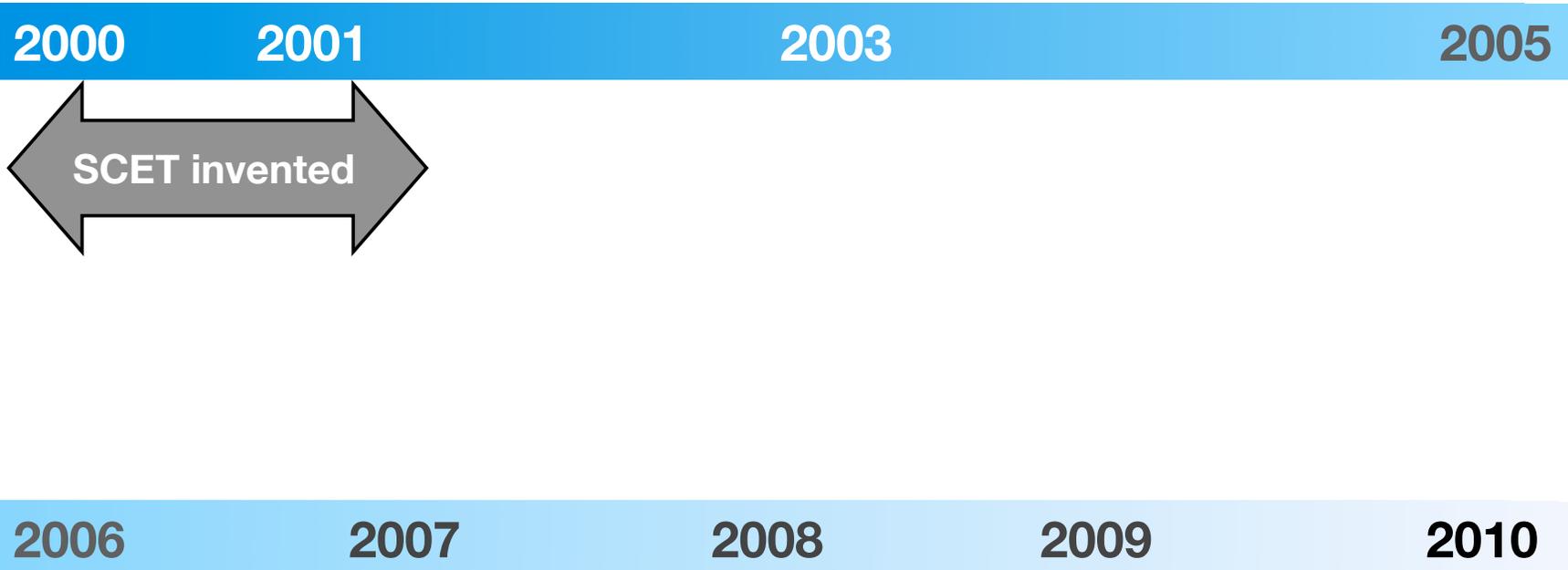
2007

2008

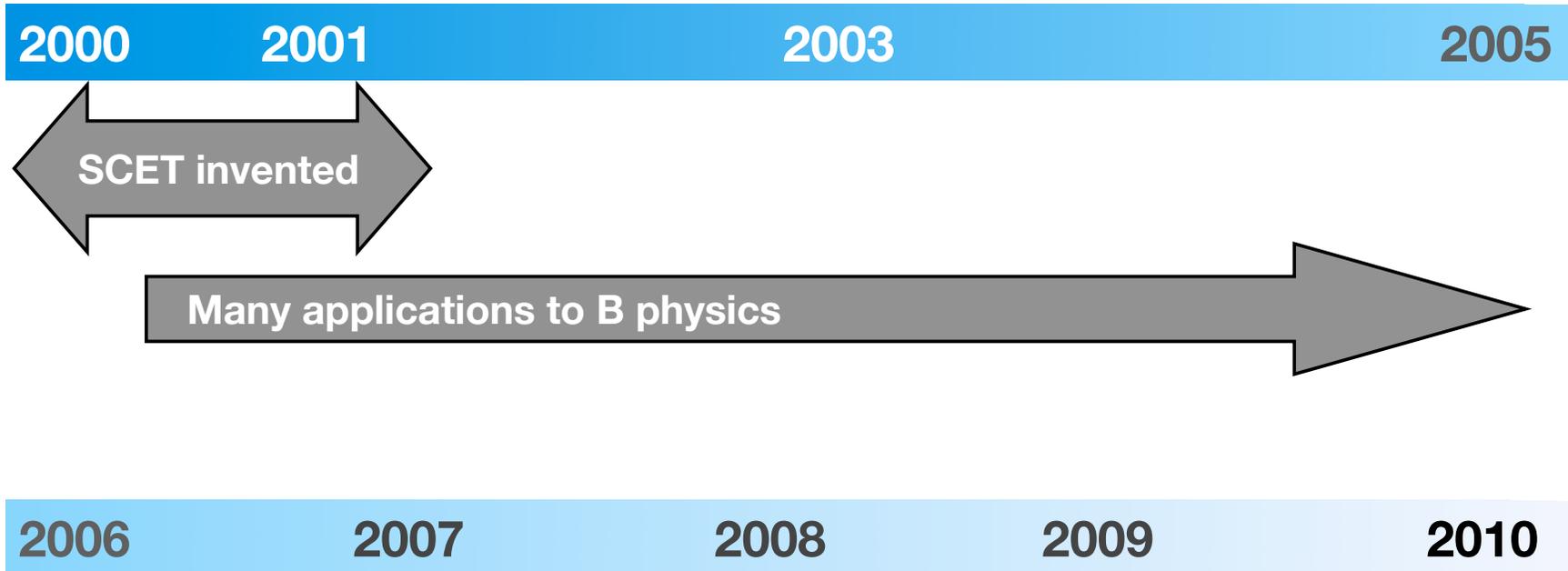
2009

2010

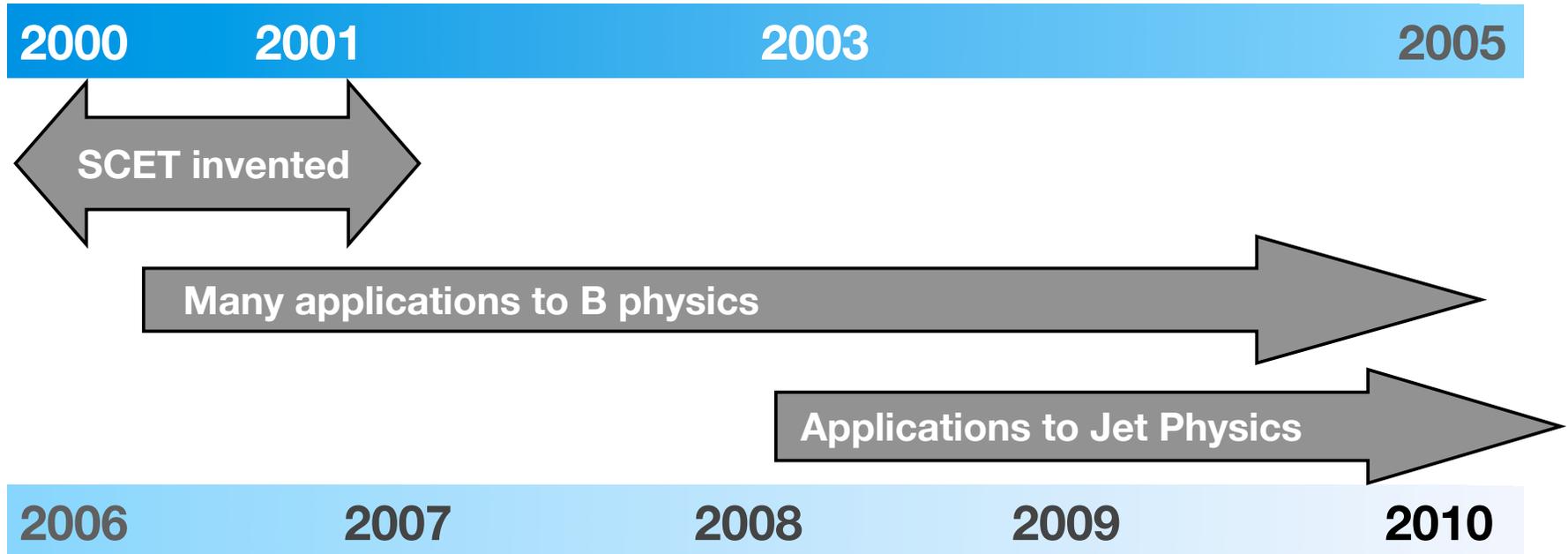
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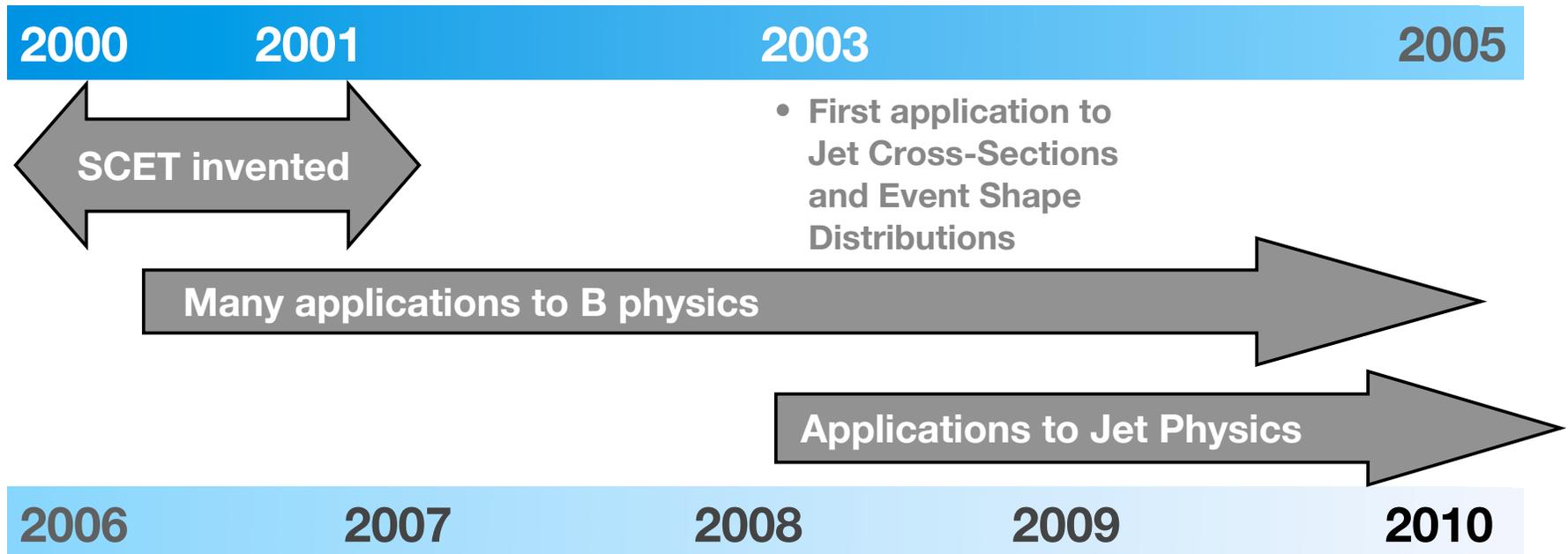
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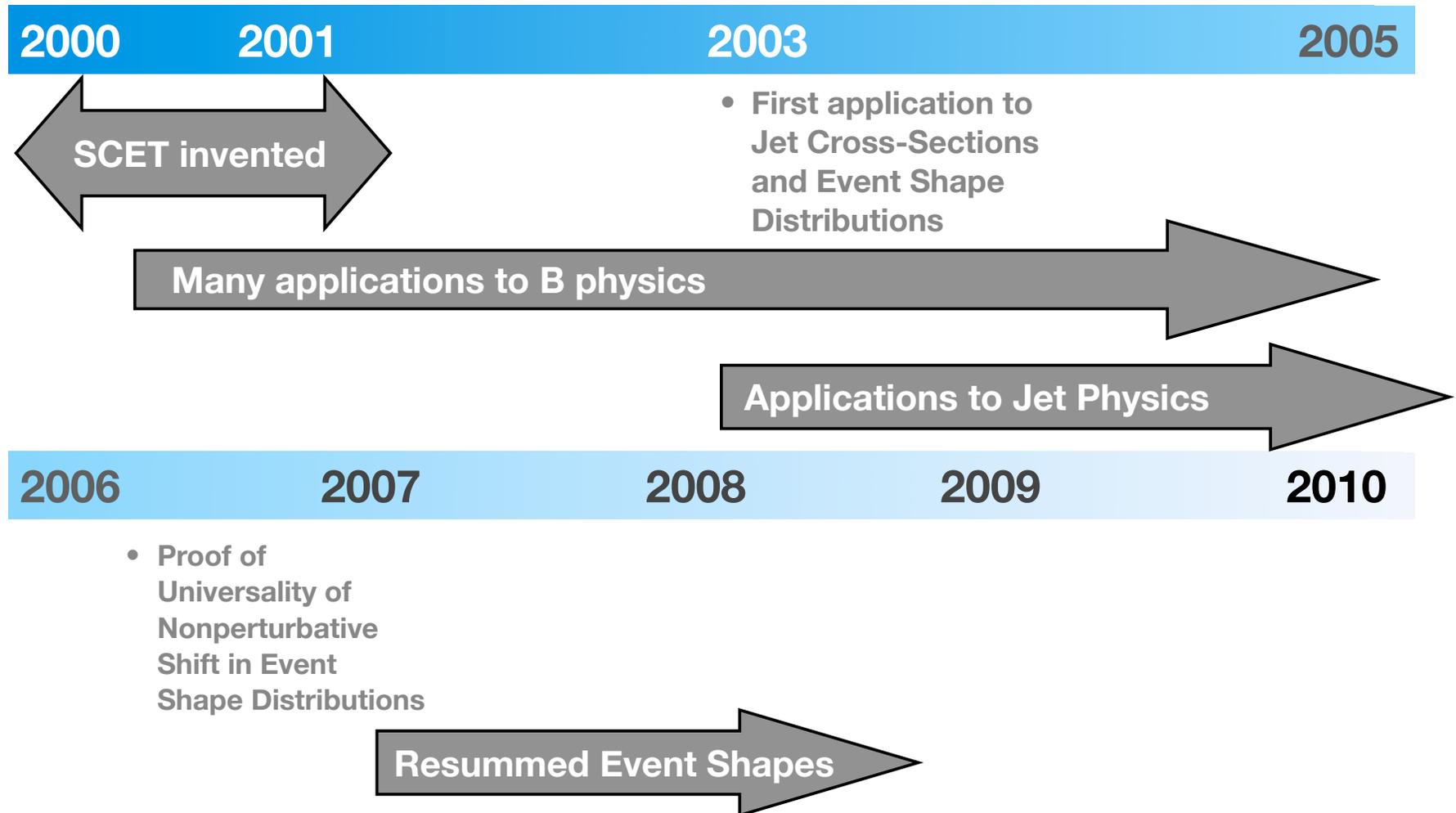
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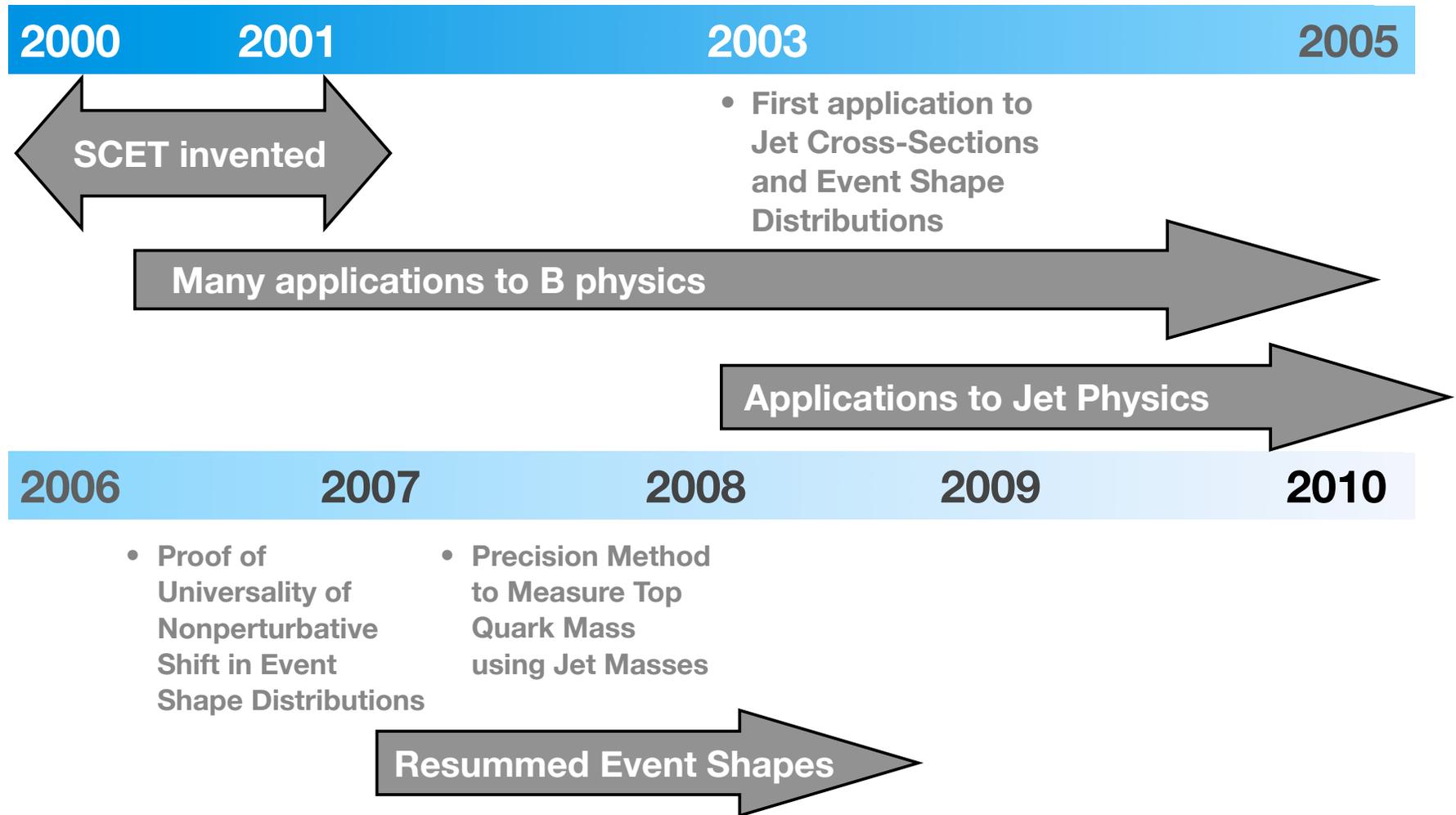
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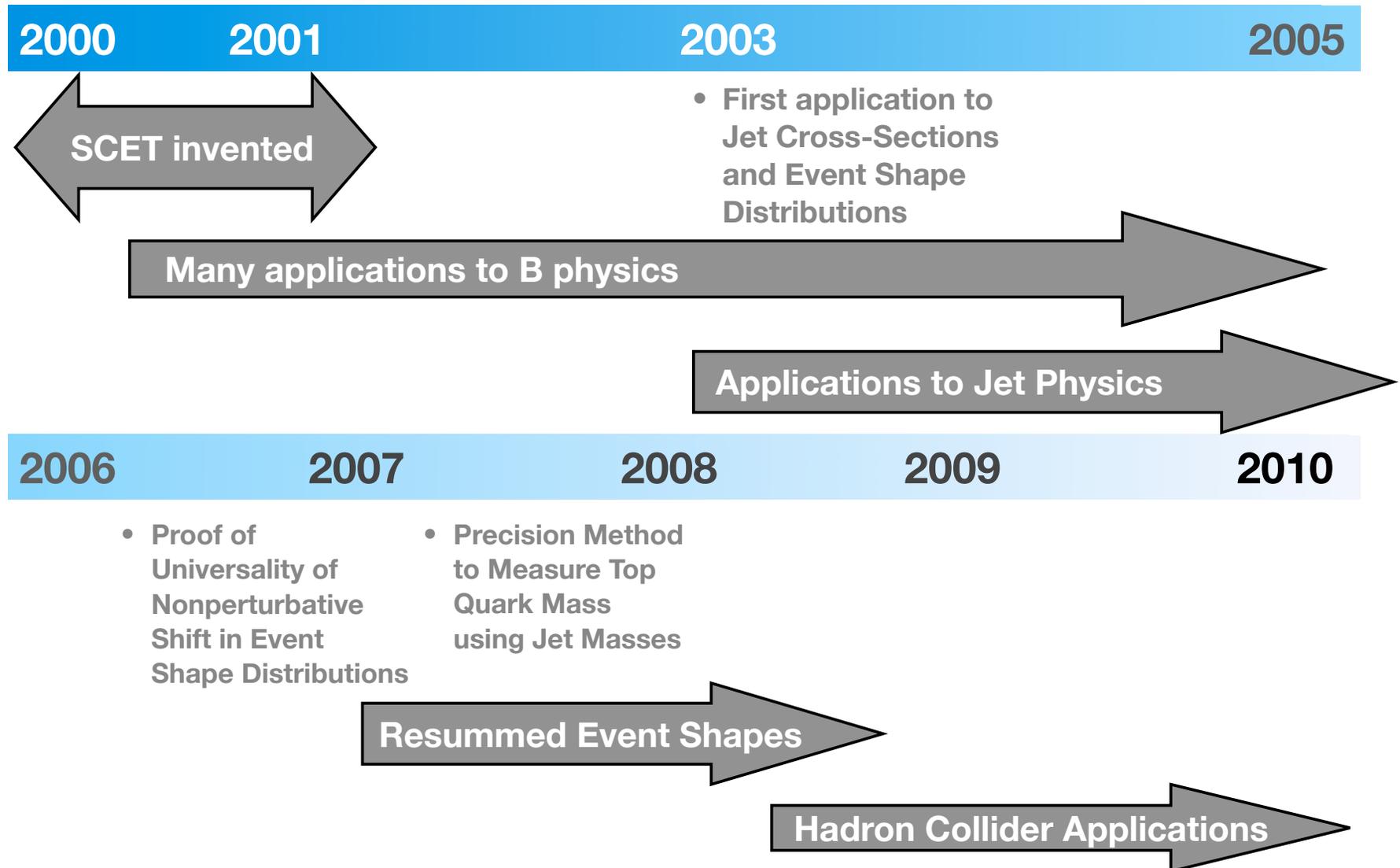
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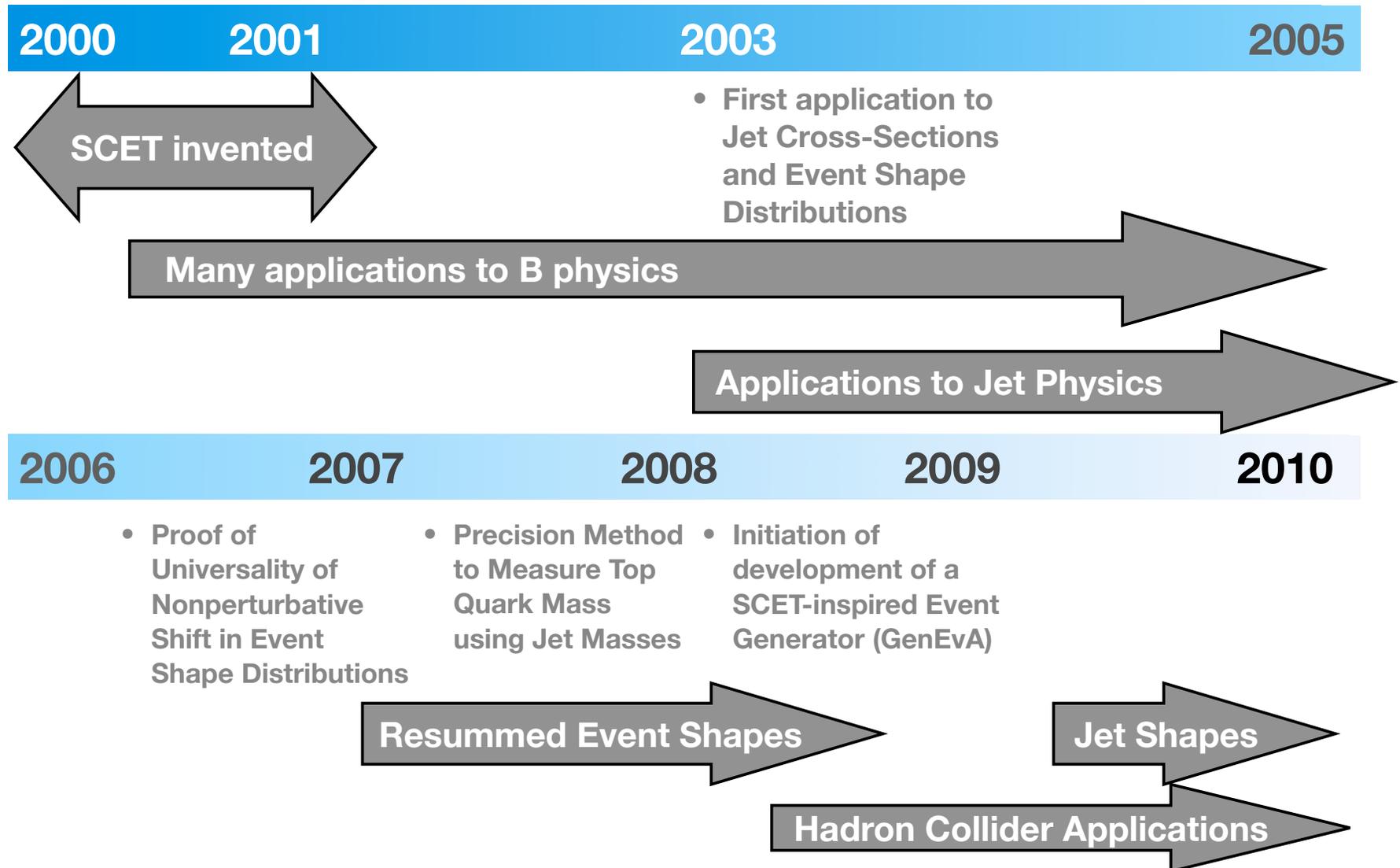
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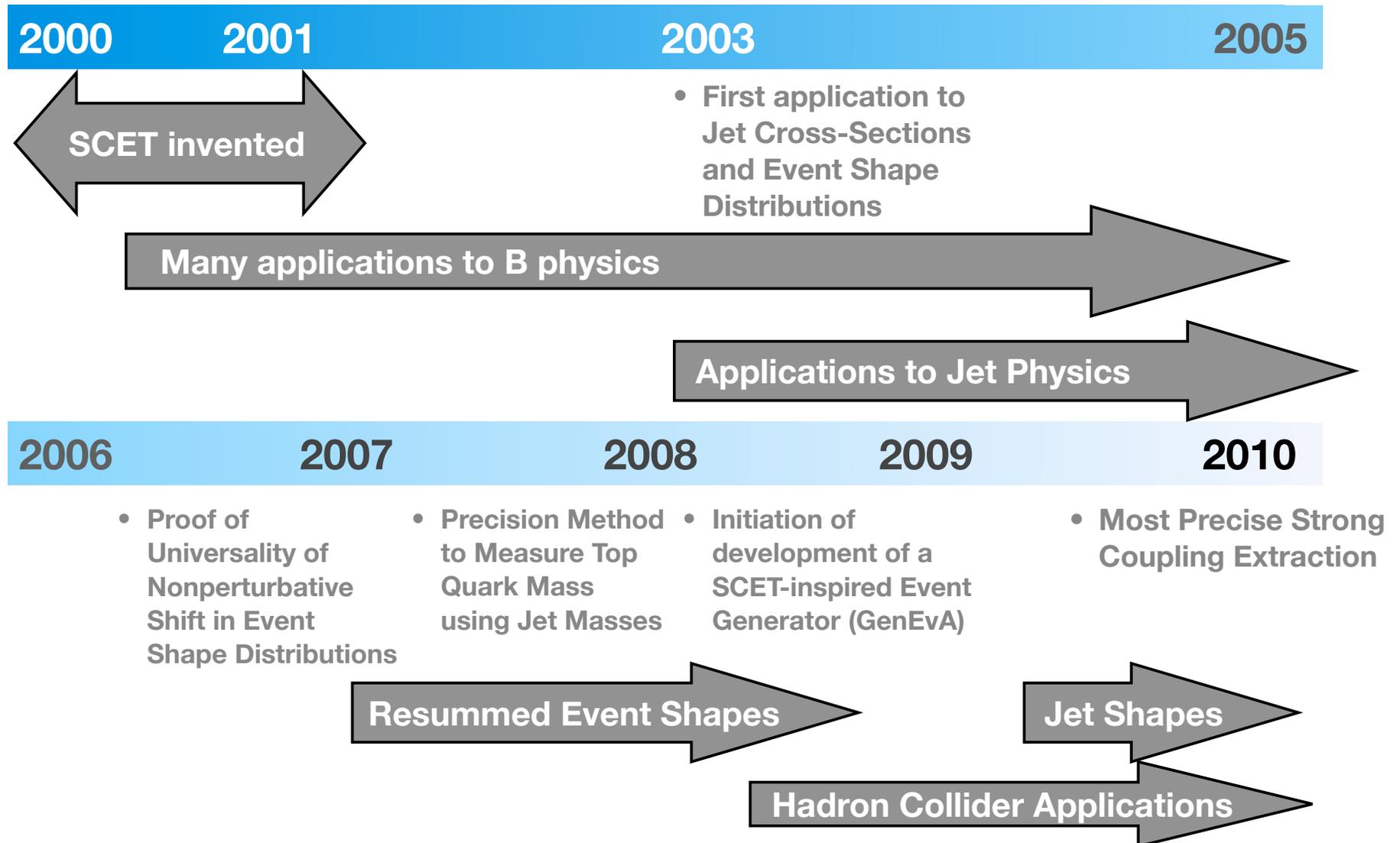
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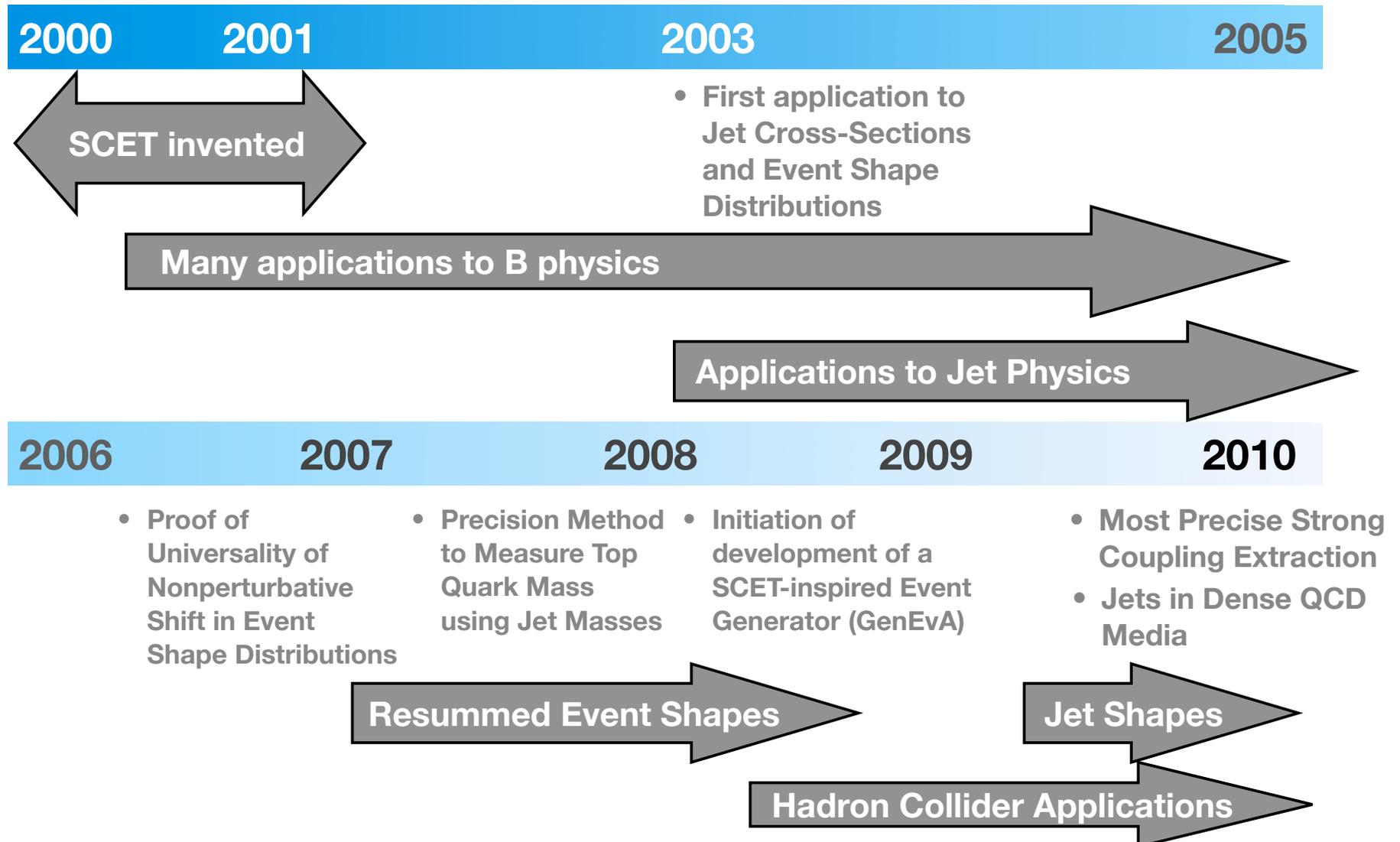
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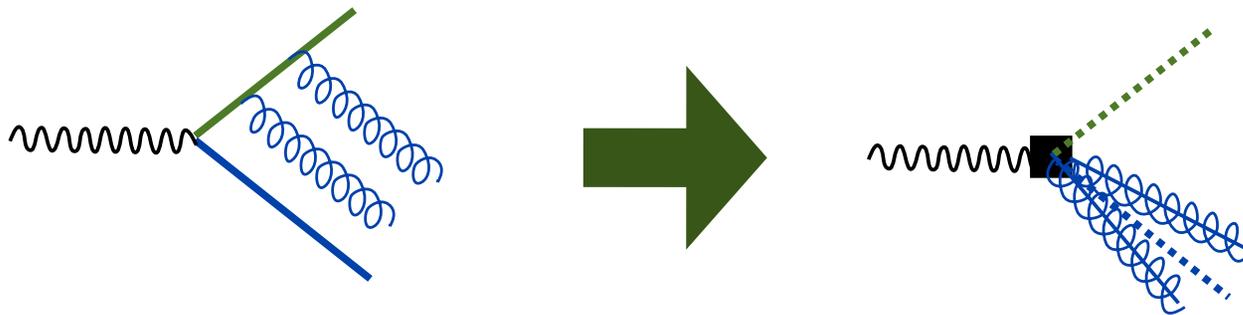


A Brief and Biased History of SCET



Going from QCD to SCET

- Integrate out hard fluctuations, e.g. when jets in other directions emit particles into the n-collinear jet



- Take eikonal limit of soft gluon attachments to collinear particles:

p k $p+k$

$= -ign^\mu T^A \frac{\not{n}}{2}$

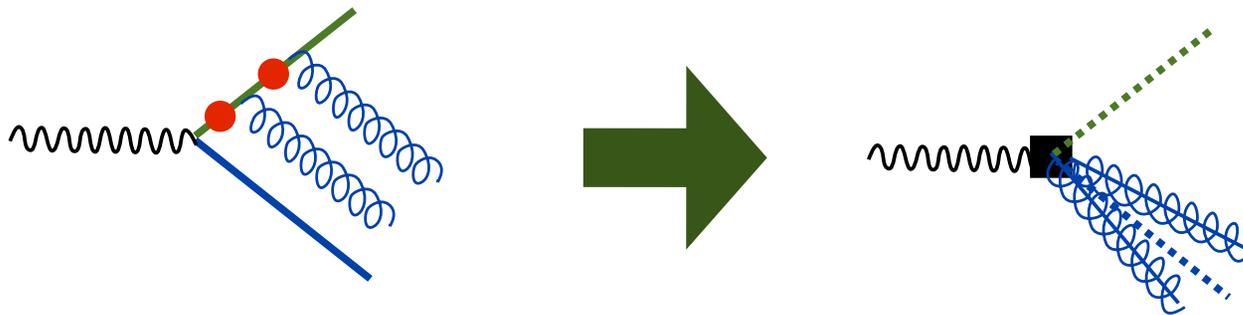
soft gluons
“see” only
direction and
color of jet

$$\frac{i(\not{p} + \not{k})}{(p+k)^2 + i\epsilon} \longrightarrow \frac{i}{n \cdot k + i\epsilon} \not{n}$$

$$Y_n(z) = P \exp \left[ig \int_0^\infty ds n \cdot A_s(ns + z) \right]$$

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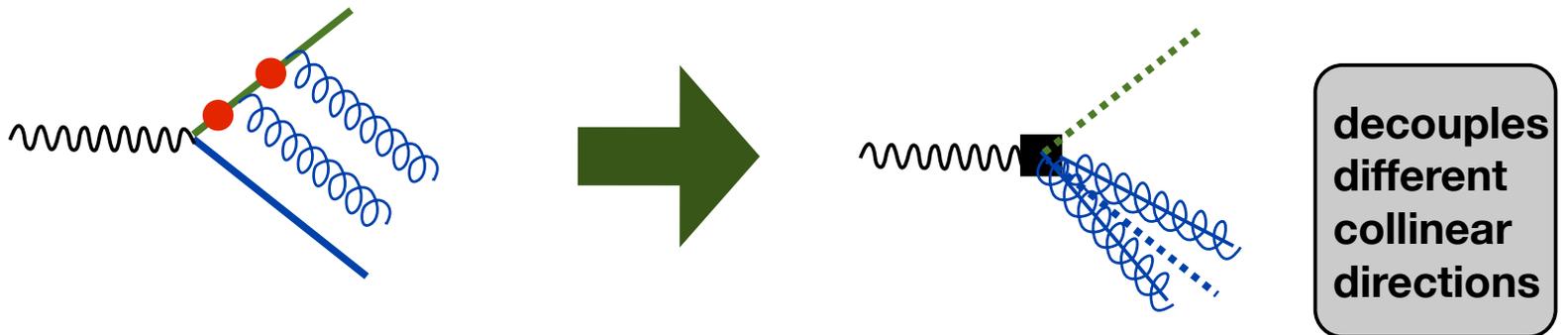
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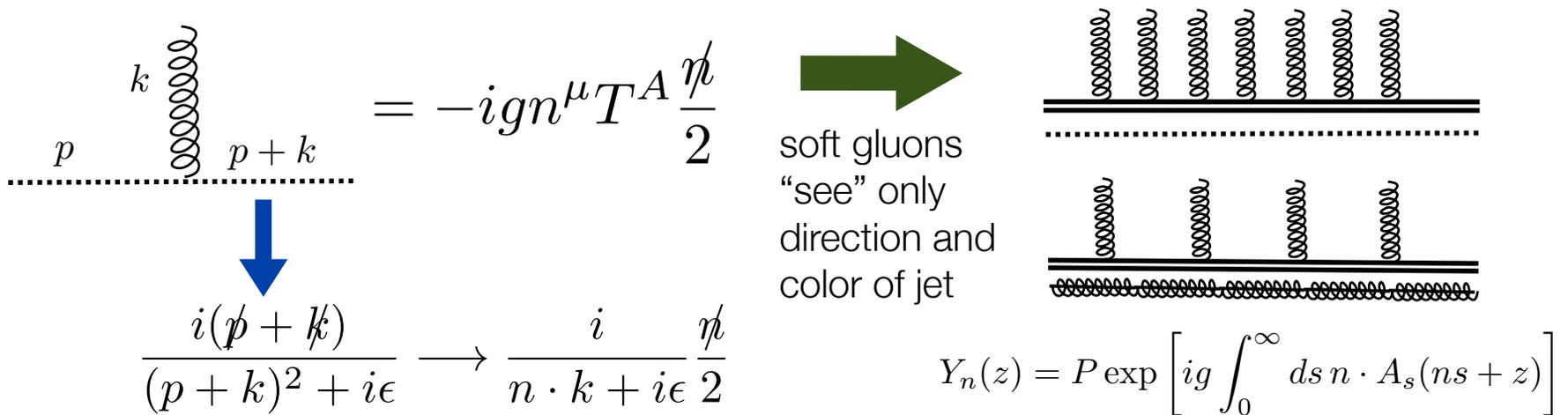
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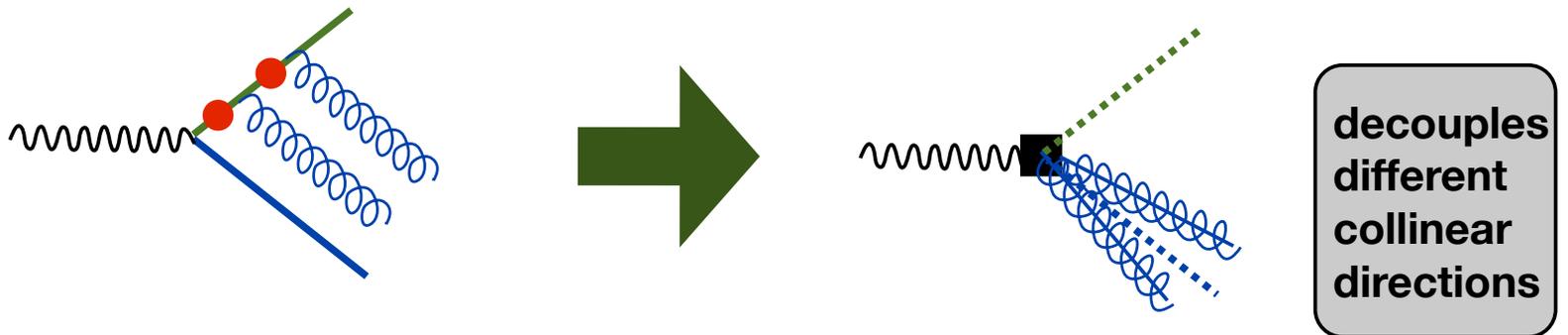


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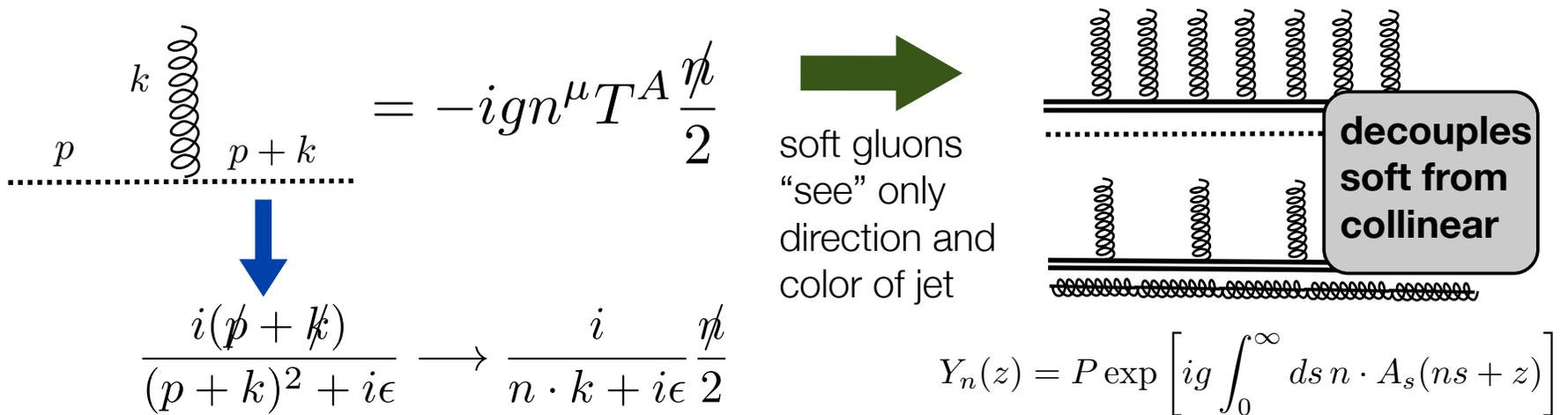


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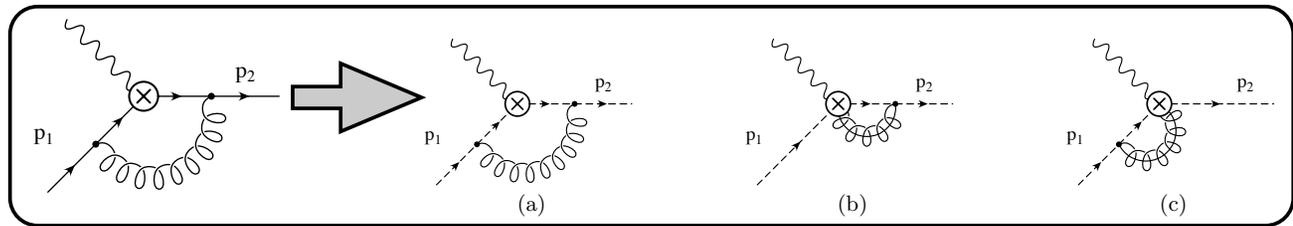
Factorization of Event Shape Distributions

Bauer, Fleming,
CL, Sterman

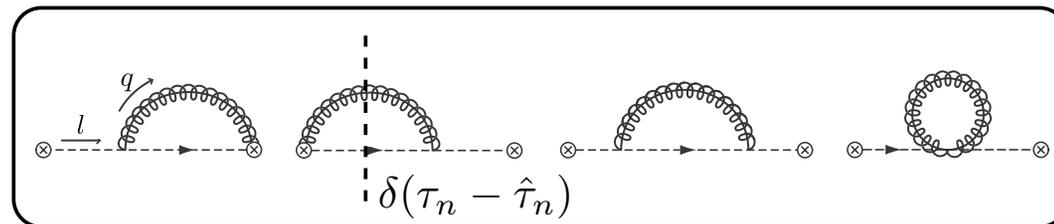
- Convolution of hard, jet, and soft functions:

$$\frac{1}{\sigma_0} \frac{d\sigma}{d\tau} = H(Q, \mu) \int d\tau_n d\tau_{\bar{n}} J_n(\tau_n; \mu) J_{\bar{n}}(\tau_{\bar{n}}; \mu) S(\tau - \tau_n - \tau_{\bar{n}}; \mu)$$

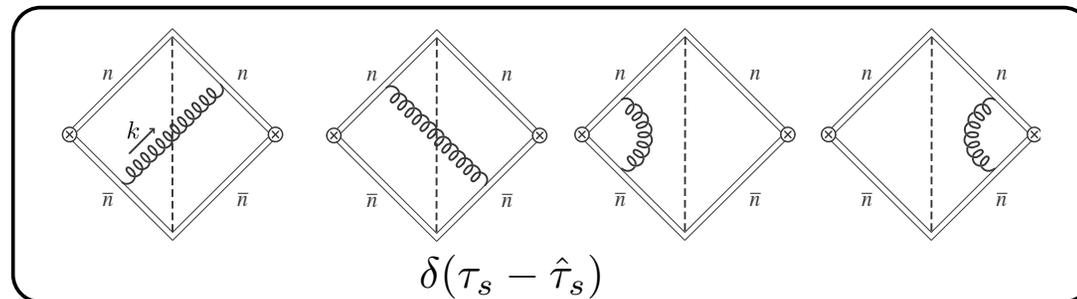
- Hard Function



- Jet Function



- Soft Function



Plan of the Talk

- Motivation: Probing QCD and New Physics with Jet Shapes
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Large Logarithms in Distributions

- Thrust distribution contains divergent terms $\alpha_s^n \frac{\ln^m \tau}{\tau}$

- or, in integrated distribution $\alpha_s^n \ln^{m+1} \tau$

$$R(\tau) = \int d\tau \frac{1}{\sigma_0} \frac{d\sigma}{d\tau}$$

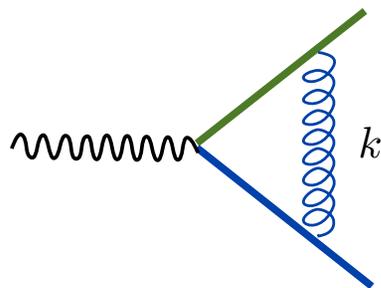
- Resum all terms up to a given order of accuracy

$$\begin{aligned}
 \ln R(\tau) \sim & \alpha_s (\ln^2 \tau + \ln \tau) \\
 & + \alpha_s^2 (\ln^3 \tau + \ln^2 \tau + \ln \tau + \dots) \\
 & + \alpha_s^3 (\ln^4 \tau + \ln^3 \tau + \ln^2 \tau + \ln \tau + \dots) \\
 & + \dots
 \end{aligned}$$

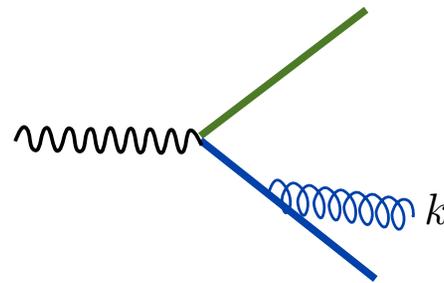
LL NLL NNLL

Where Do Logs Come from?

- Infrared divergences in **virtual** loop graphs canceled by soft and collinear divergences in **real** emission diagrams integrated over all phase space:

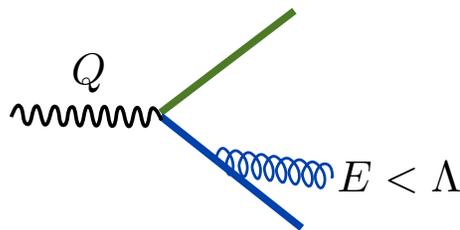


$$\int d\Pi_2 \int d^4k \Rightarrow \frac{A}{\epsilon_{\text{IR}}}$$

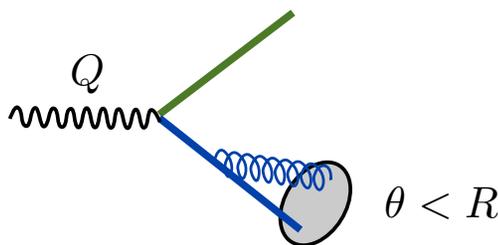


$$\int d\Pi_3 \Rightarrow -\frac{A}{\epsilon_{\text{IR}}}$$

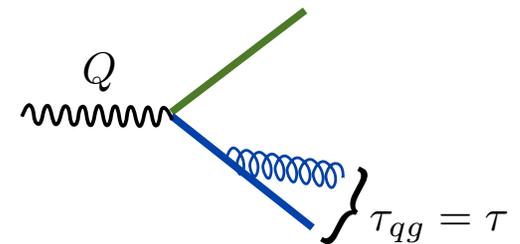
- Restrictions on real emission phase space leave the cancellation **incomplete** and generate **large logarithms**:



$$\int_{E < \Lambda} d\Pi_3 \Rightarrow -\frac{A}{\epsilon_{\text{IR}}} + \ln \frac{\Lambda}{Q}$$



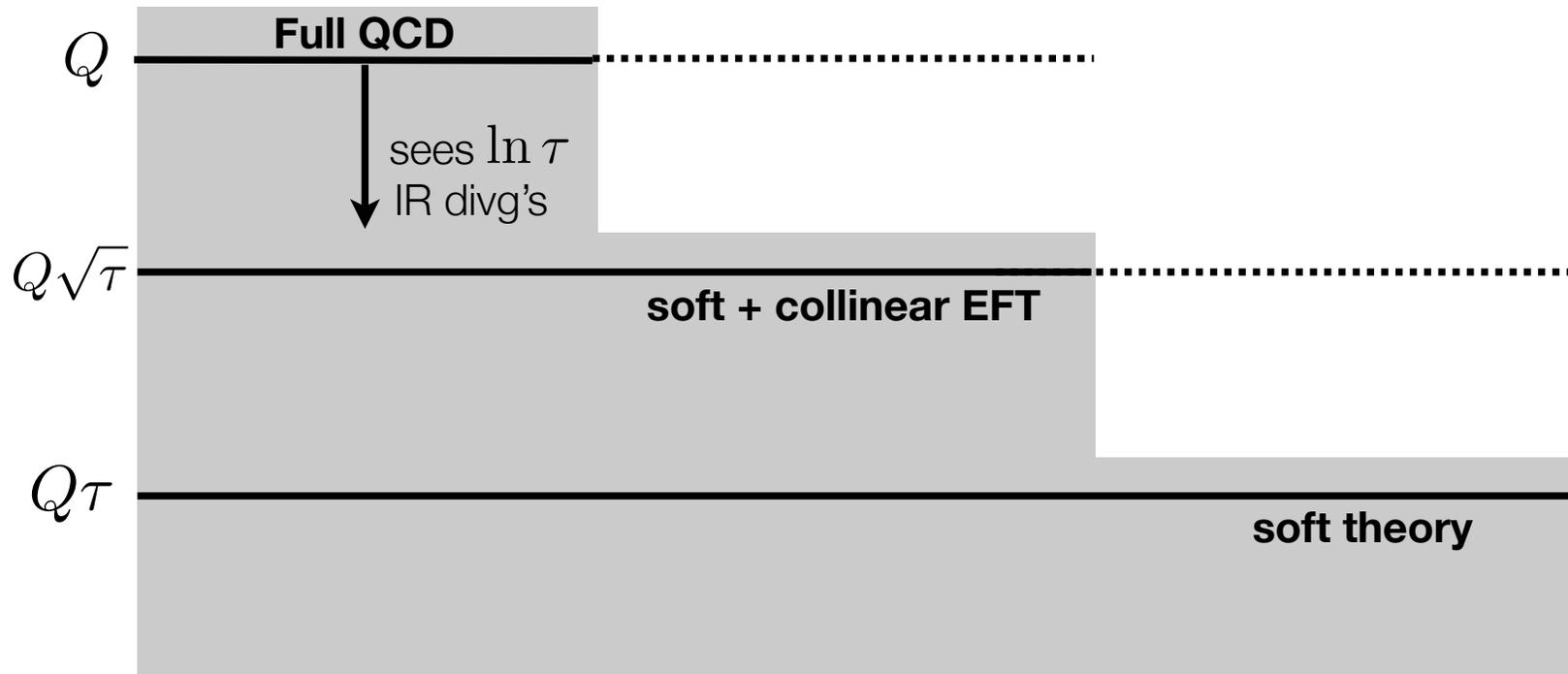
$$\int_{\theta < R} d\Pi_3 \Rightarrow -\frac{A}{\epsilon_{\text{IR}}} + \ln \tan \frac{R}{2}$$



$$\int d\Pi_3 \delta(\tau_{qg} - \tau) \Rightarrow -\frac{A}{\epsilon_{\text{IR}}} + \ln \tau$$

IR Resummation from UV Evolution

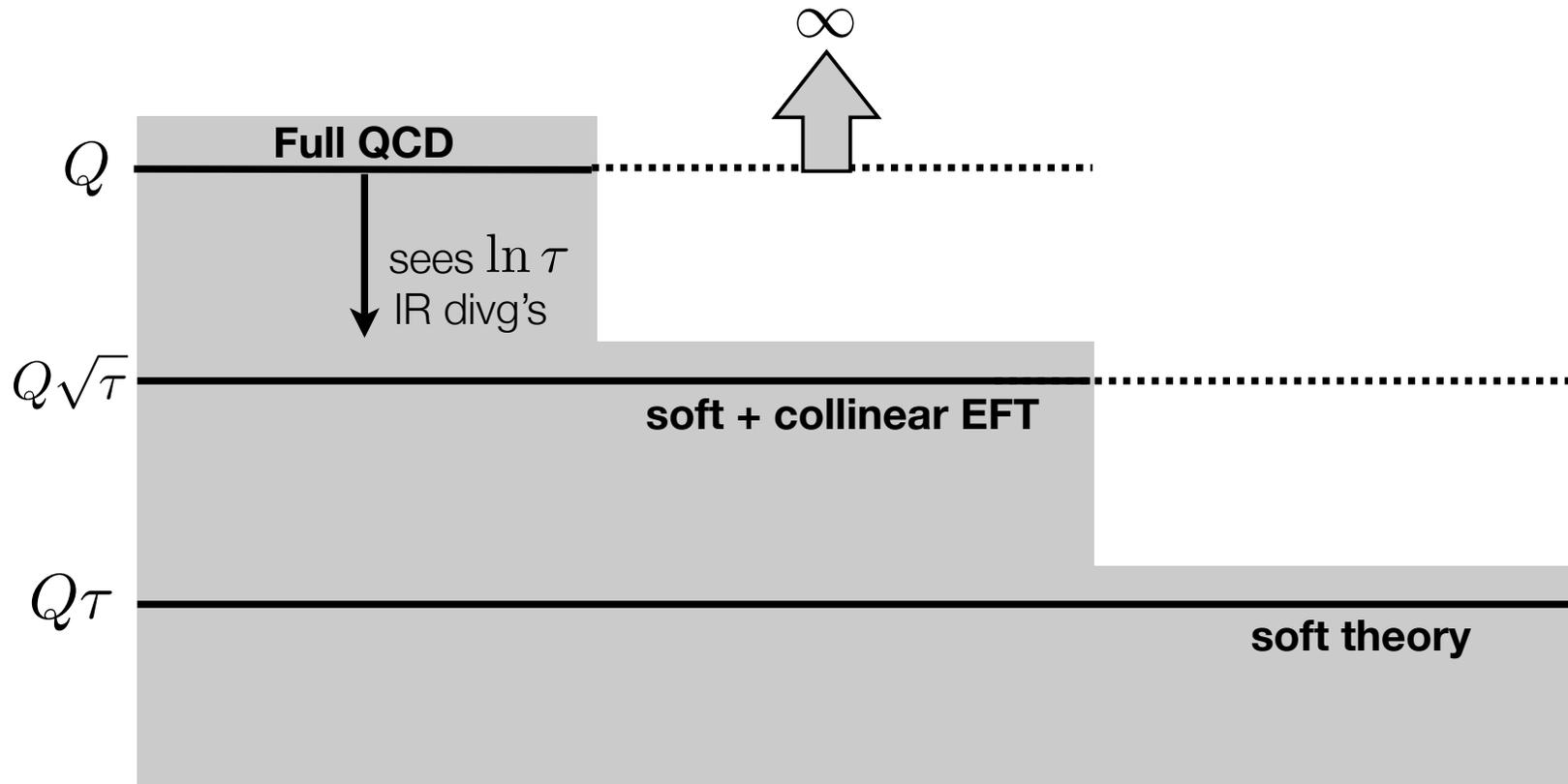
- IR divergences in the full theory are converted to UV divergences in the effective theory, e.g. for thrust distribution:



- Logs can now be resummed using renormalization group evolution with the factorization scale μ

IR Resummation from UV Evolution

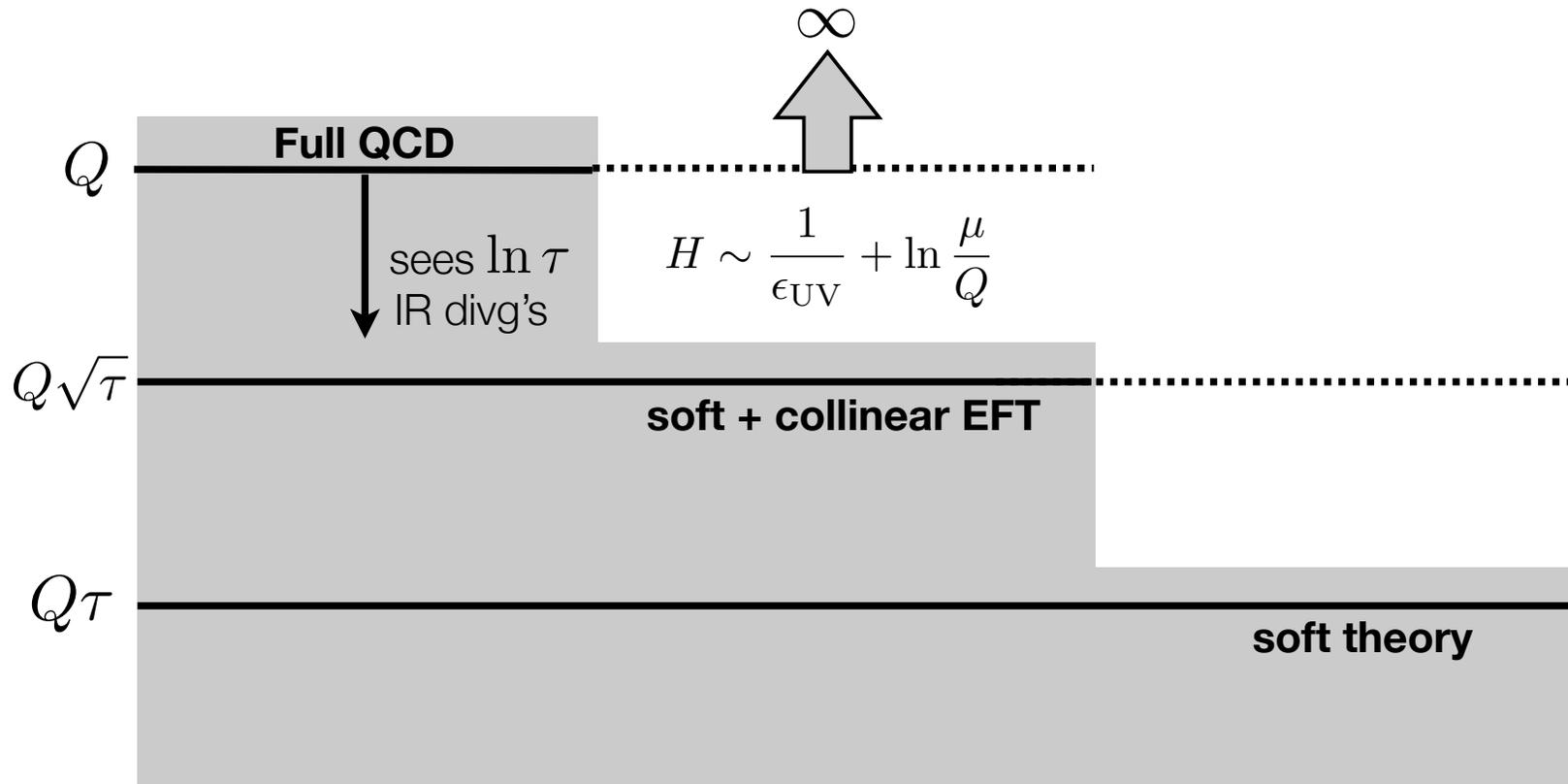
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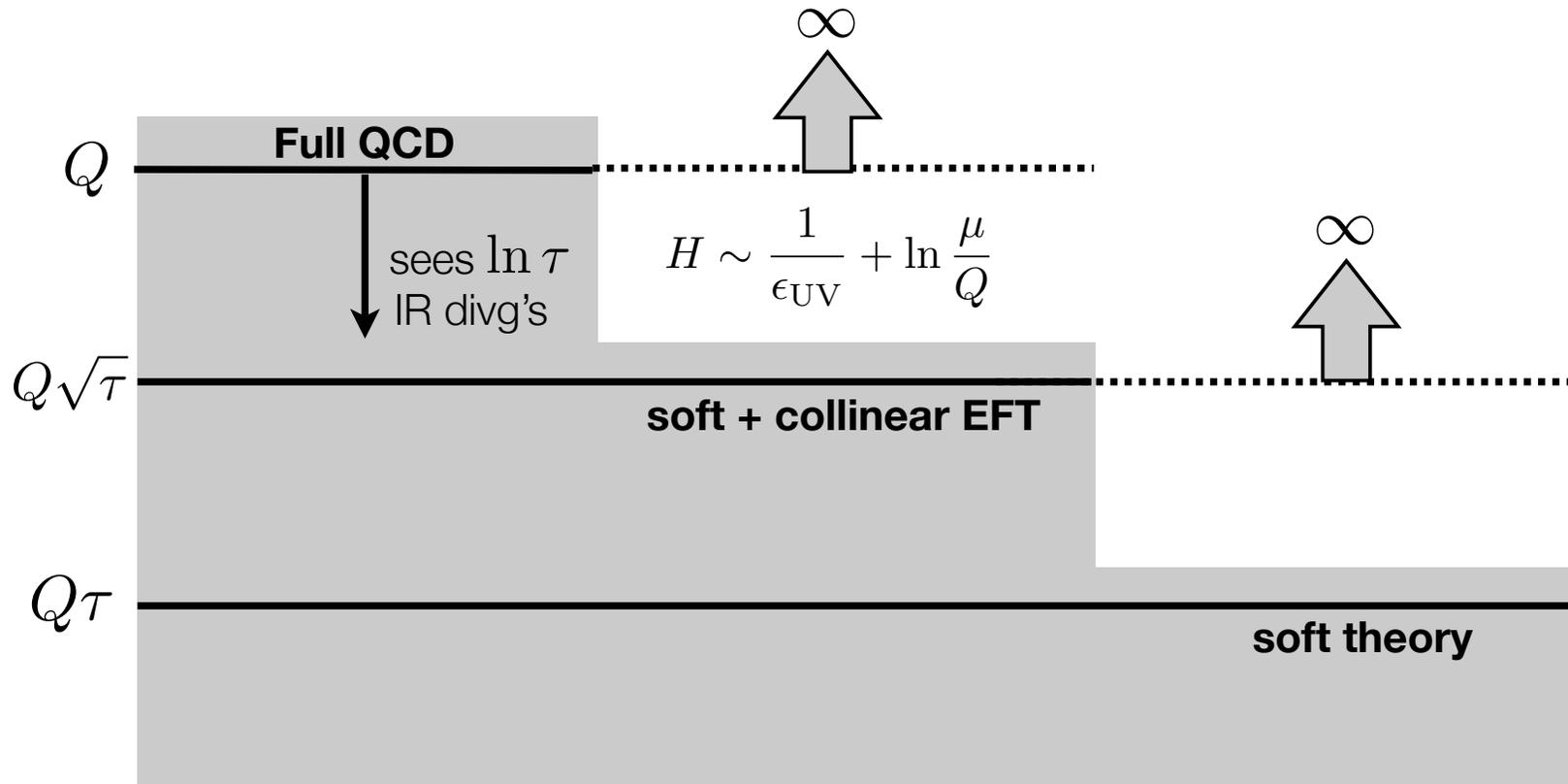
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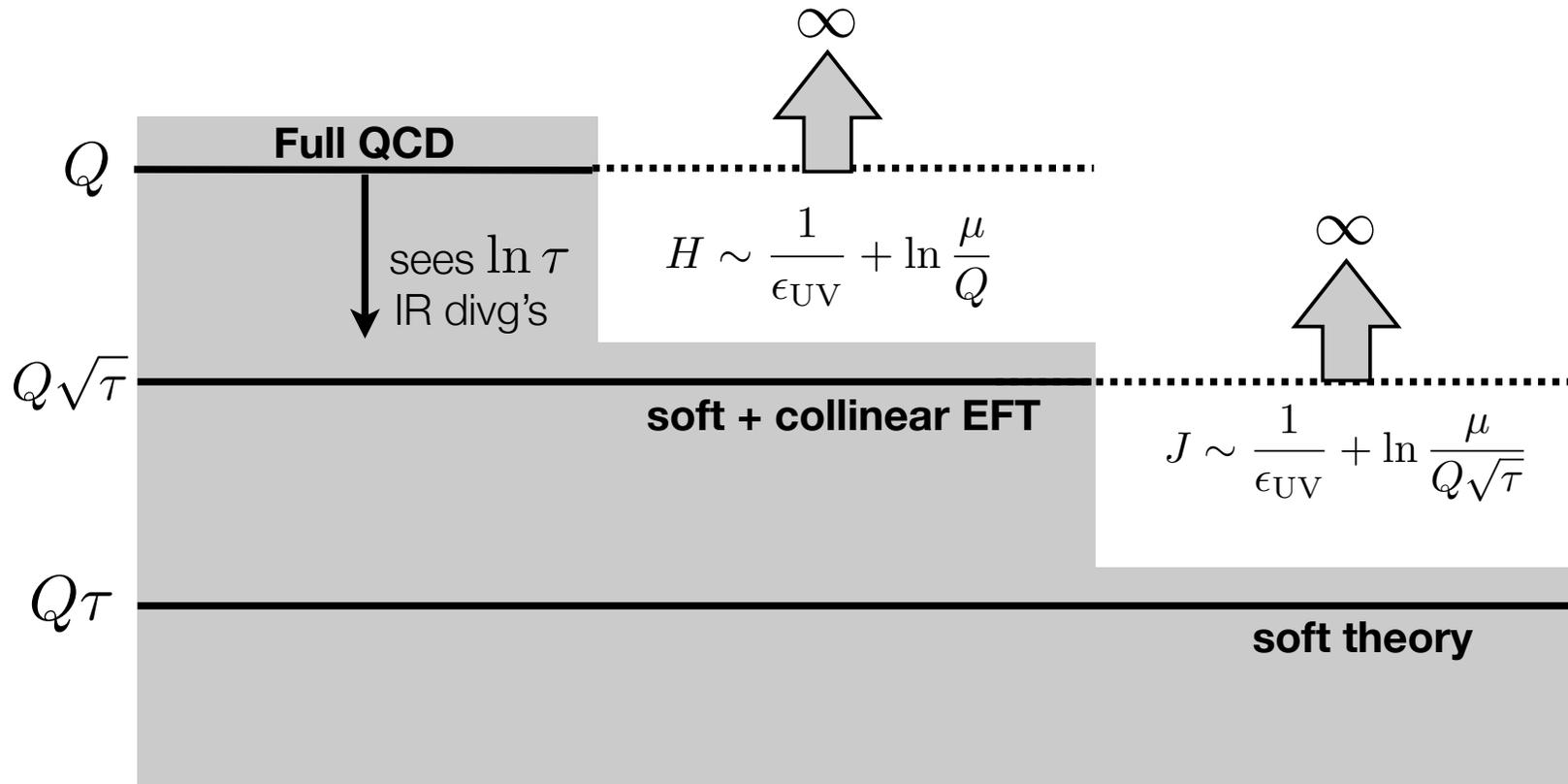
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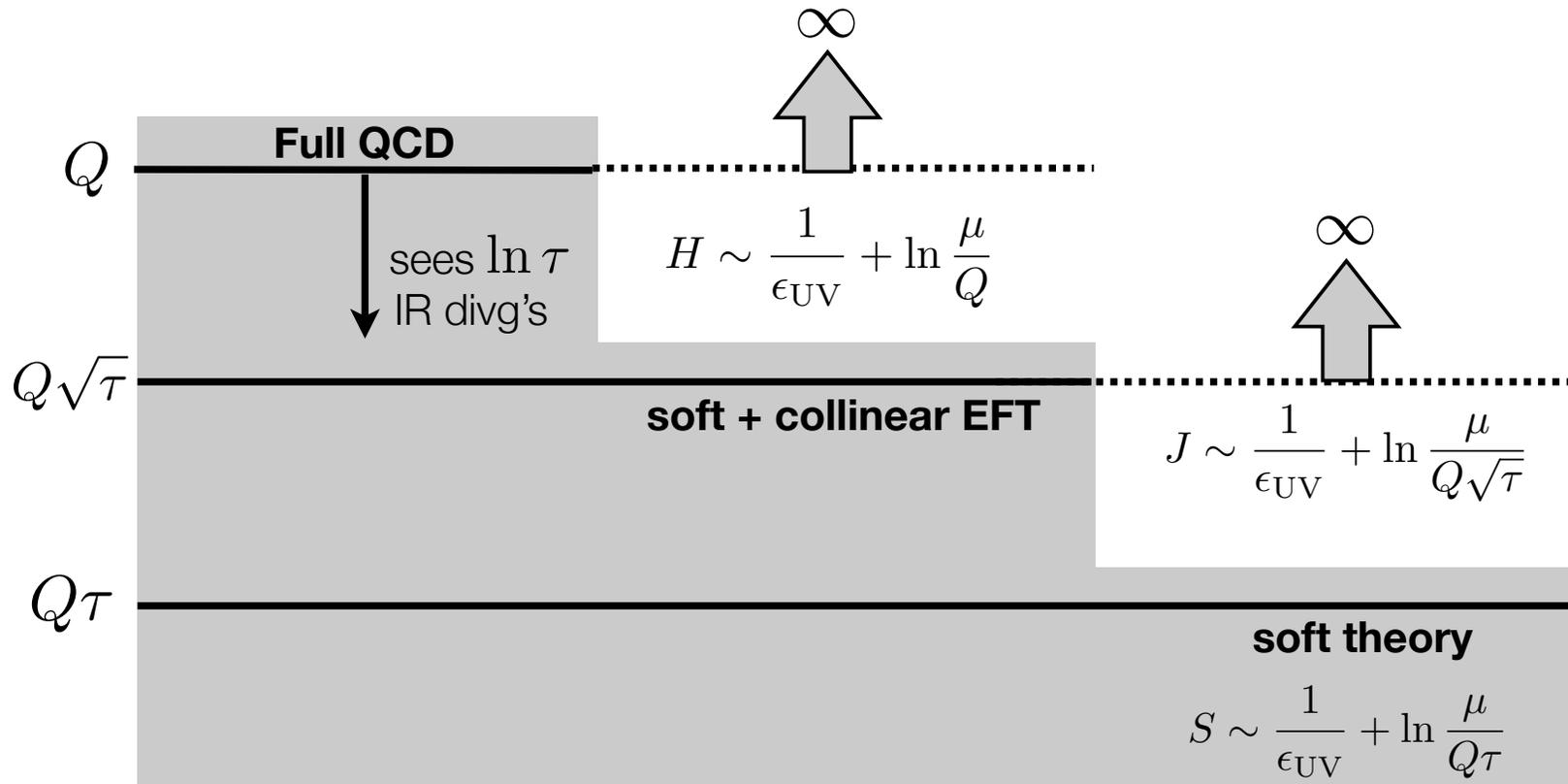
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Universal Soft Power Corrections

- Soft power corrections shift mean values of event shapes

$$\langle e \rangle = \langle e \rangle_{\text{PT}} + c_e \frac{\mathcal{A}}{Q}$$

conjecture from single soft gluon emission:
Dokshitzer, Webber
(1995, 1997)

c_e observable dependent,
calculable coefficient

\mathcal{A} universal nonperturbative parameter

- Stronger version: simple shift of full distribution

$$\frac{d\sigma}{de}(e) = \frac{d\sigma}{de} \left(e - c_e \frac{\mathcal{A}}{Q} \right) \Big|_{\text{PT}}$$

- SCET: First rigorous proof (and **field theory** definition of \mathcal{A}) from boost invariance of soft Wilson lines in factorization theorem
- In general, use NP soft shape function to smear out perturbative distribution, with universal information for first moment

Universal Soft Power Corrections

- Soft power corrections shift mean values of event shapes

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conjecture from single soft gluon emission:
Dokshitzer, Webber (1995, 1997)

c_e observable dependent, calculable coefficient

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Use factorization to prove true

CL, Sterman (2007)

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In general not true

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Angularities: a general class of event shapes

- Generalization of thrust:

Berger, Kucs, Sterman

$$\tau_a = \frac{1}{Q} \sum_i E_i (\sin \theta_i)^a (1 - |\cos \theta_i|)^{1-a} = \frac{1}{Q} \sum_i |\mathbf{p}_i^T| e^{-|\eta_i|(1-a)}$$

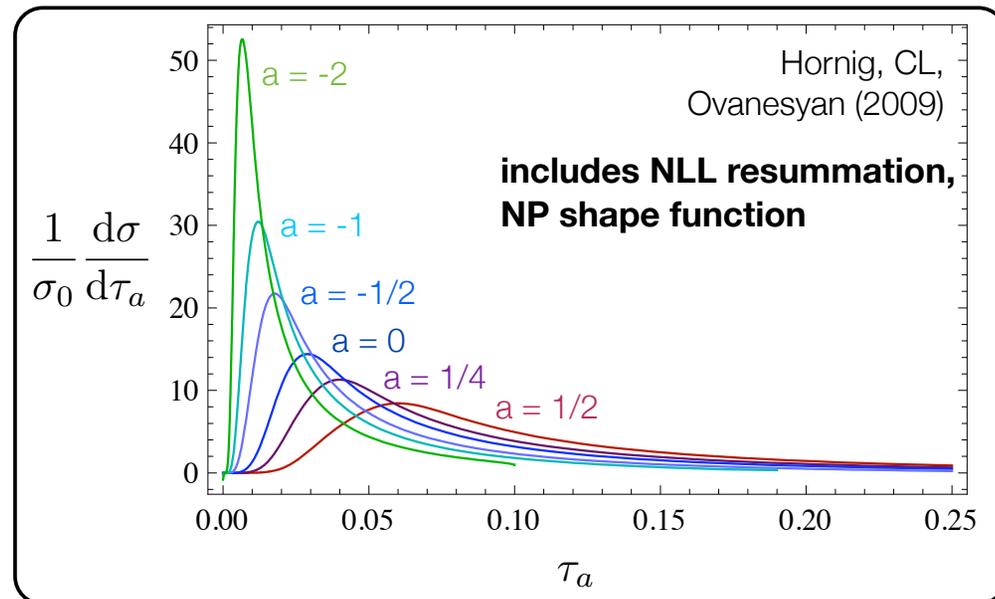
$a = 0$ thrust

$a = 1$ broadening

infrared safety: $-\infty < a < 2$

factorizability: $-\infty < a < 1$

Peak region dominated
by narrower jets for large a ,
wider jets for small a



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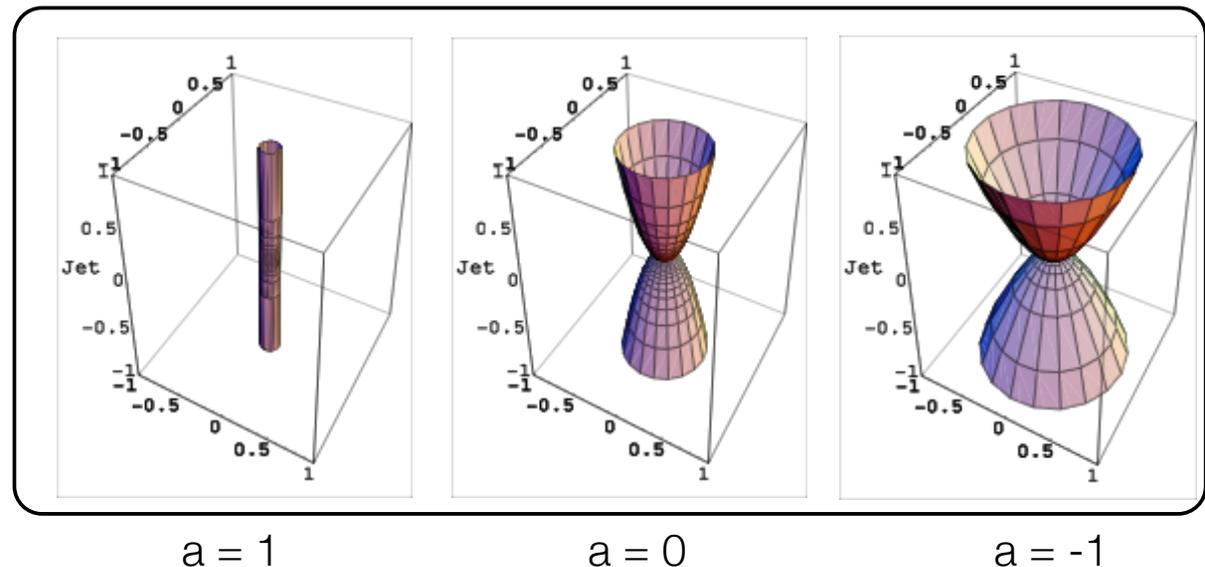
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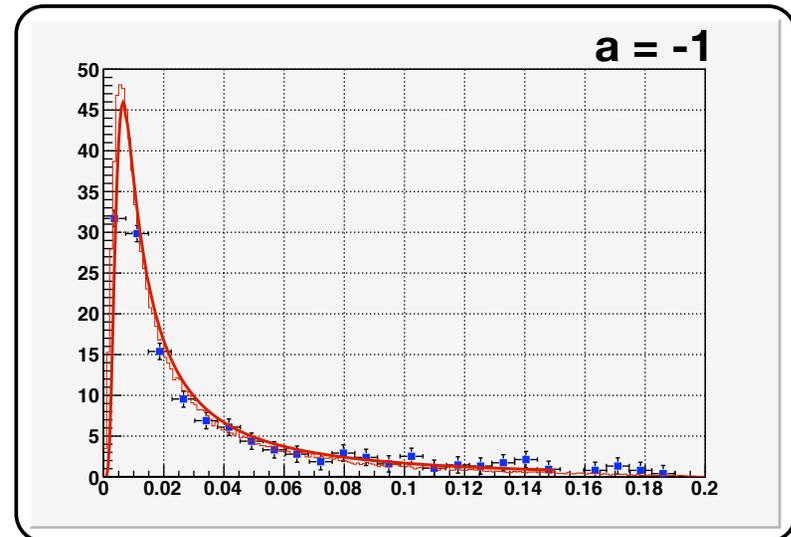
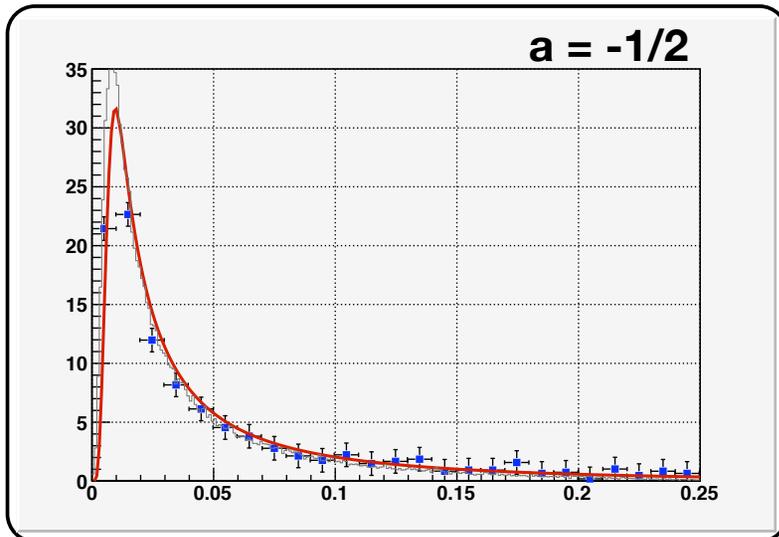
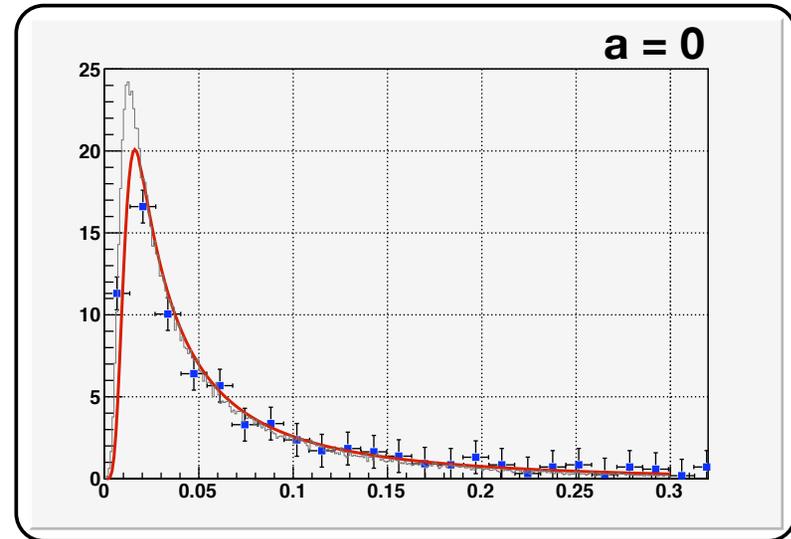
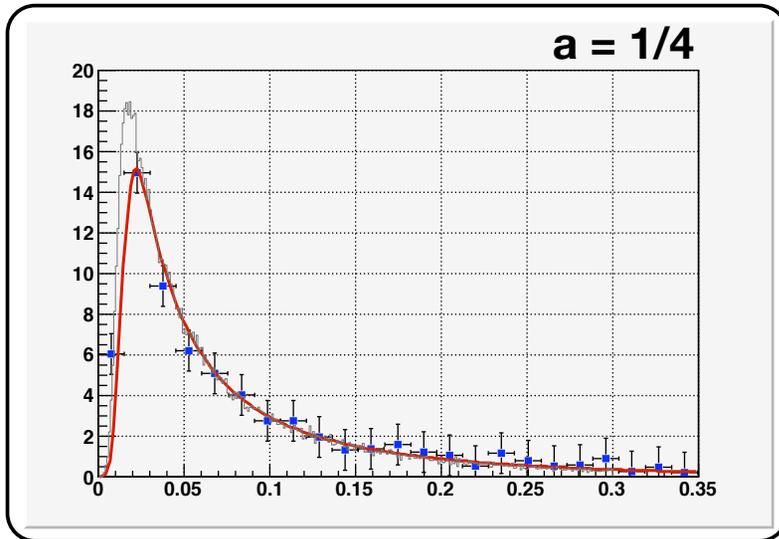
Peak region dominated by narrower jets for large a , wider jets for small a



Angularity Distributions at $Q = 197 \text{ GeV}$:

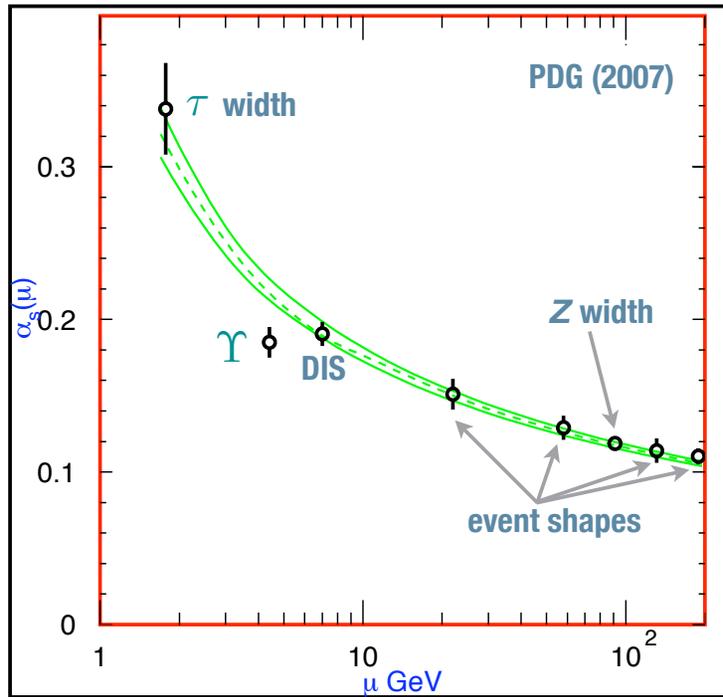
NLL Theory vs. **Experiment** (and PYTHIA8)

Th: Hornig, CL, Ovanesyan
Ex: P. Jindal (LEP), PhD Thesis



Precision extractions of strong coupling

$$\alpha_s(M_Z) = 0.1176 \pm 0.0020 \text{ (PDG)}$$

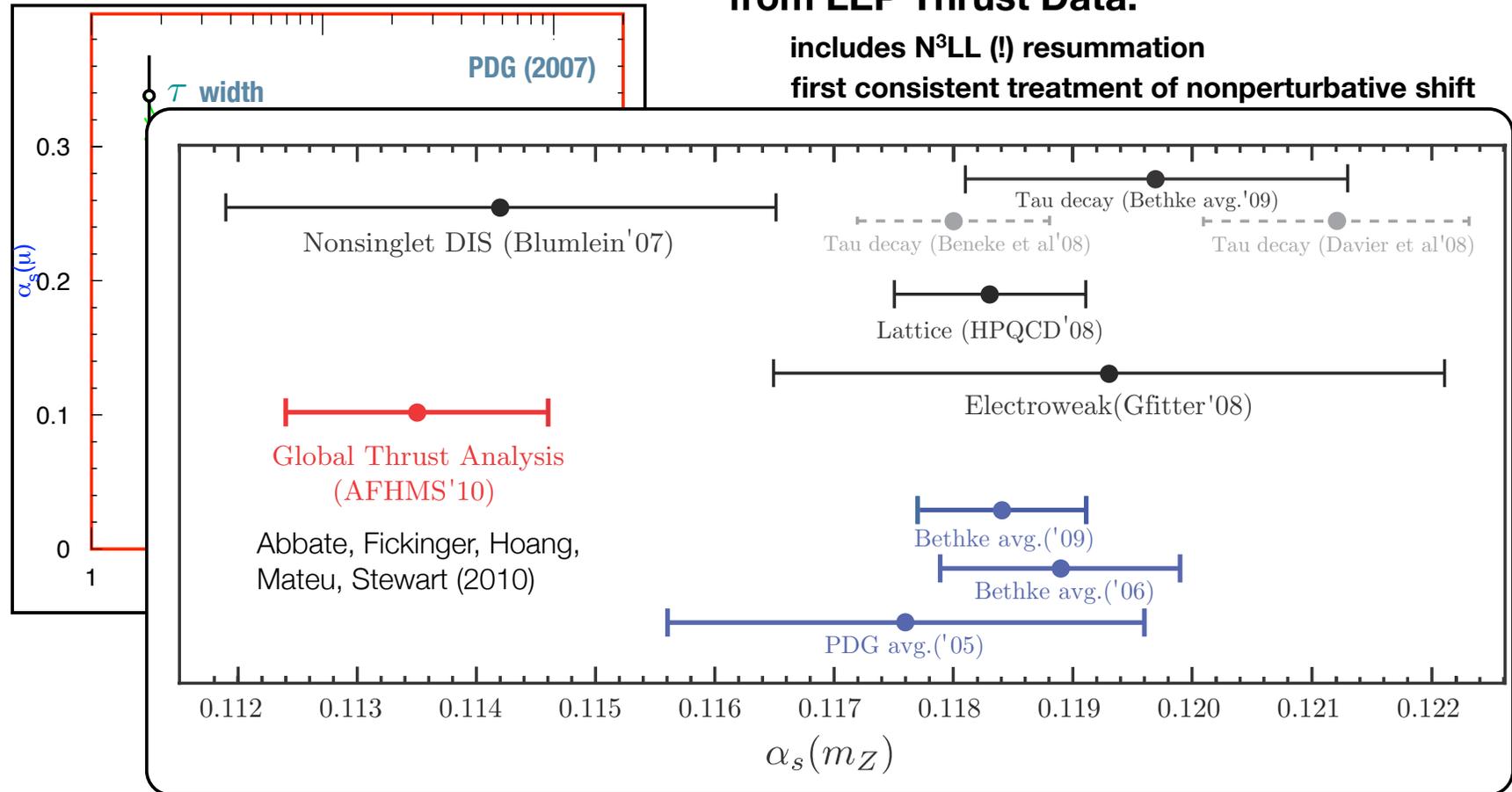


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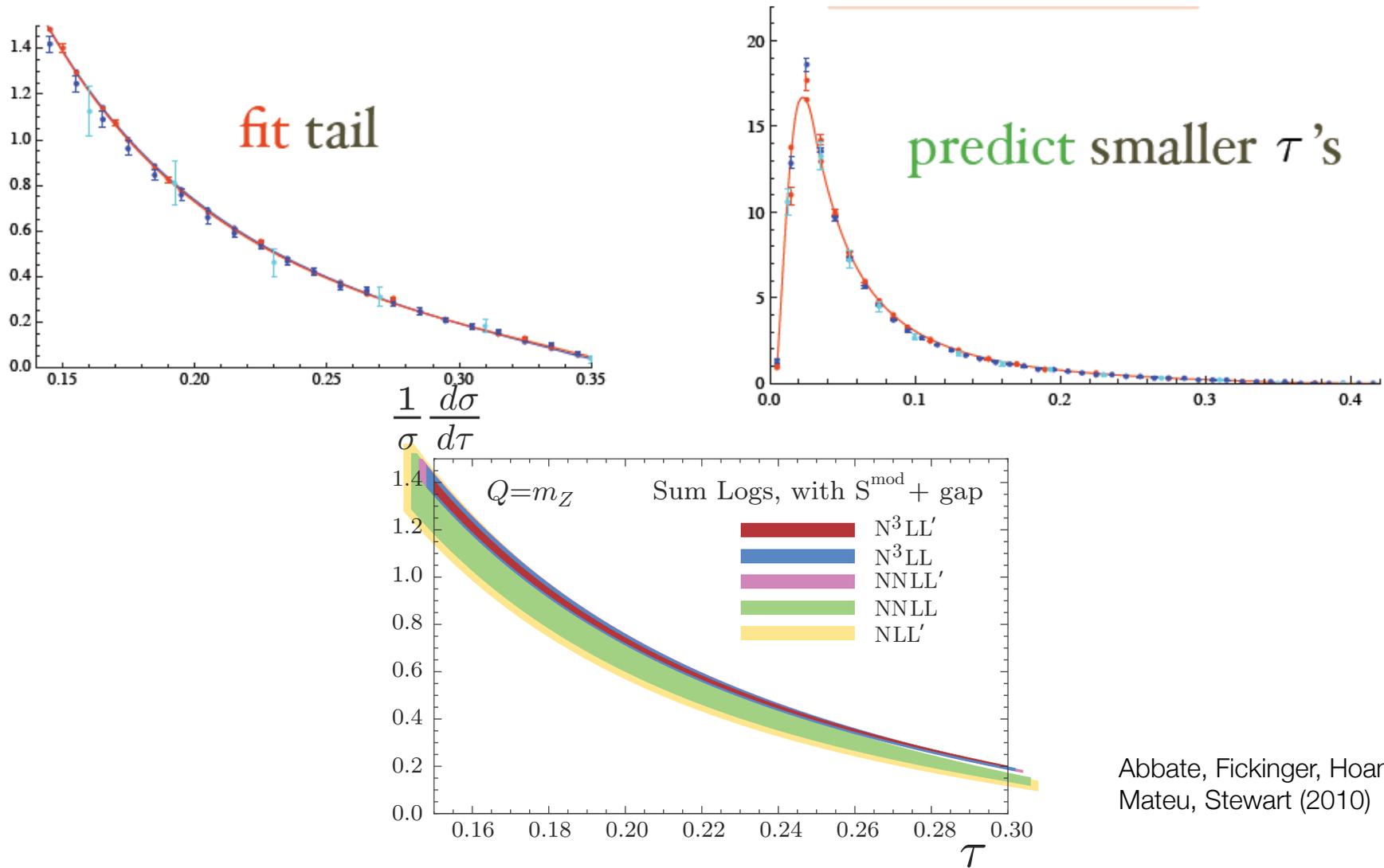
New SCET-based Extraction from LEP Thrust Data:

includes N³LL (!) resummation
 first consistent treatment of nonperturbative shift



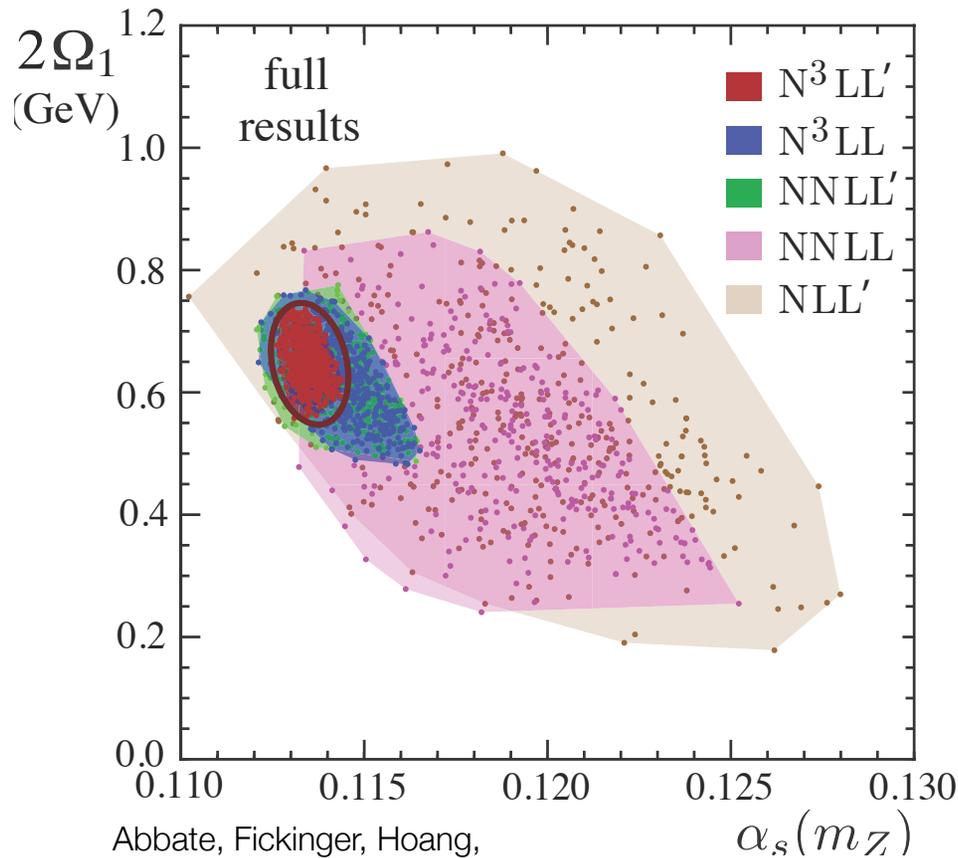
$$\alpha_s(M_Z) = 0.1135 \pm (0.0002)_{\text{exp}} \pm (0.0005)_{\text{had}} \pm (0.0009)_{\text{pert}}$$

Precision extractions of strong coupling

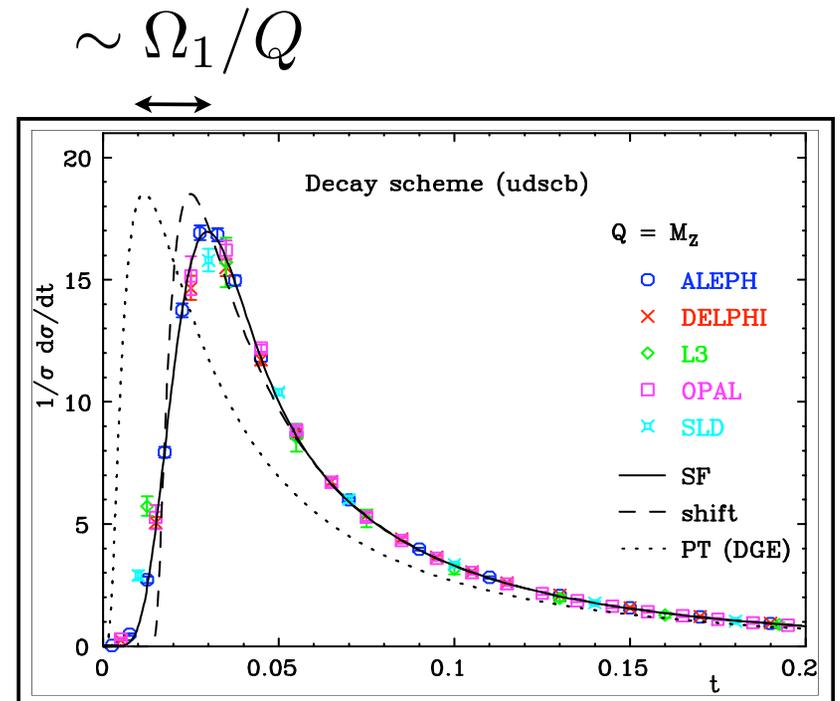


Abbate, Fickinger, Hoang, Mateu, Stewart (2010)

Precision extractions of strong coupling



Abbate, Fickinger, Hoang,
Mateu, Stewart (2010)



$$\bar{\Omega}_1 \equiv \frac{1}{2N_c} \langle 0 | \text{tr} \bar{Y}_{\bar{n}}(0) Y_n(0) i\partial_\tau Y_n^\dagger(0) \bar{Y}_{\bar{n}}^\dagger(0) | 0 \rangle \quad (\overline{MS})$$

$$i\partial_\tau \equiv \theta(i\bar{n} \cdot \partial - i\bar{n} \cdot \partial) i\mathbf{n} \cdot \partial + \theta(i\mathbf{n} \cdot \partial - i\bar{n} \cdot \partial) i\bar{n} \cdot \partial$$

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Strategies to Probe Jet Substructure

- **“Jet Grooming”**: Filtering, Trimming, Pruning...

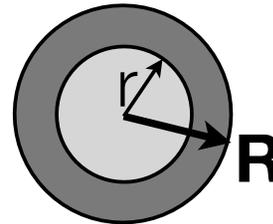
- well-suited for implementation in Monte Carlo or expt.

Butterworth, Davison, Rubin, Salam; Thaler, Wang; Ellis, Vermilion, Walsh; Kaplan, Rehermann, Schwartz, Tweedie; Plehn, Salam, Spannowsky; Kribs, Martin, Roy, Spannowsky; ...



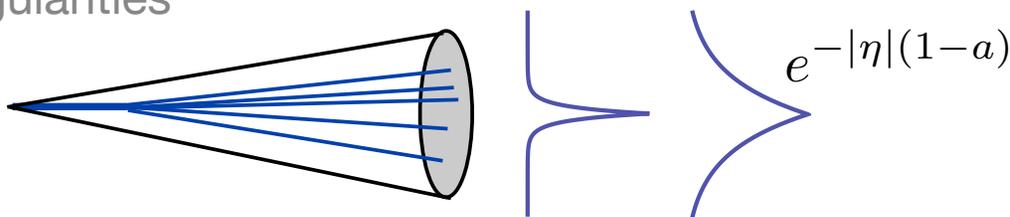
- **Jet Shapes**

- “The” Jet Shape $\Psi(r/R)$



Ellis, Kunzst, Soper;
CDF

- Angularities



Berger, Kucs, Sterman

Almeida, Lee, Perez,
Sterman, Sung, Virzi

$$\tau_a(\text{jet}) = \frac{1}{2E_{\text{jet}}} \sum_{i \in \text{jet}} |\mathbf{p}_i^T| e^{-|\eta_i|(1-a)}$$

Ellis, Hornig, CL, Vermilion,
Walsh

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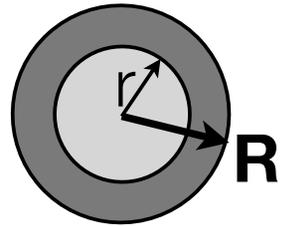
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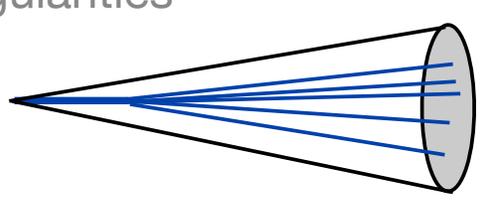
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$$\Psi(r/R)$$



Ellis, Kunzst, Soper;
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$$e^{-|\eta|(1-a)}$$

Berger, Kucs, Sterman

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Ellis, Hornig, CL, Vermilion,
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$$\tau_a(\text{jet}) = \frac{1}{2E_{\text{jet}}} \sum_{i \in \text{jet}} |\mathbf{p}_i^T| e^{-|\eta_i|(1-a)}$$

Well suited for analytical calculations

From Event Shapes to Jet Shapes

- Modify notion of event shape to the “jet shape” of a single jet:

$$\tau_a(\text{jet}) = \frac{1}{2E_{\text{jet}}} \sum_{i \in \text{jet}} |\mathbf{p}_i^T| e^{-|\eta_i|(1-a)}$$

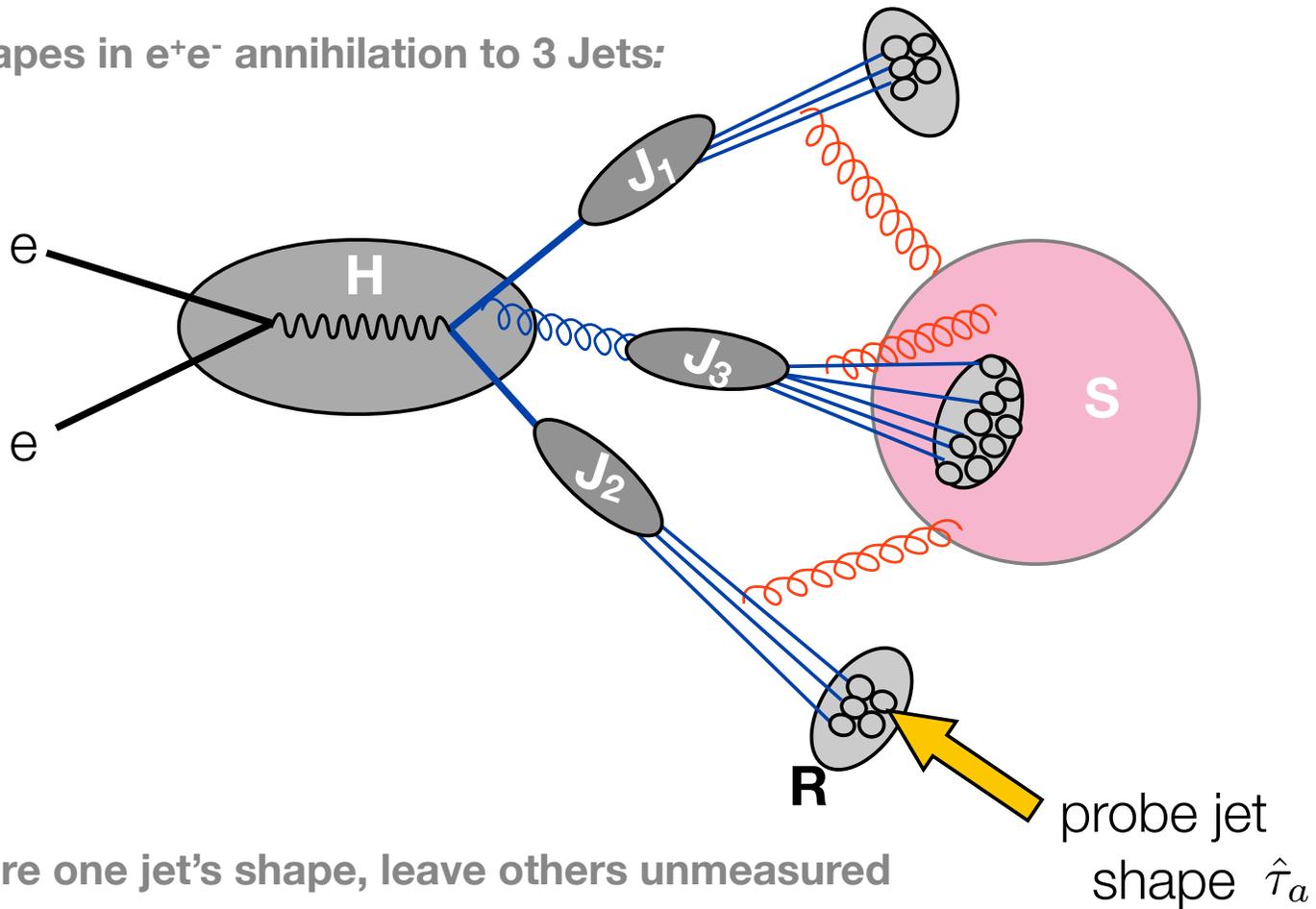
↓
(theoretical convenience)

- Requires specification of jet definition (algorithm)
- Transverse momentum and rapidity with respect to jet axis, identified by algorithm
- Requires new formulation of factorization theorem, new calculation of hard, jet, and soft functions

Ellis, Hornig, CL, Vermilion, Walsh
(2009, 2010)

Factorization of Jet Shape Distributions

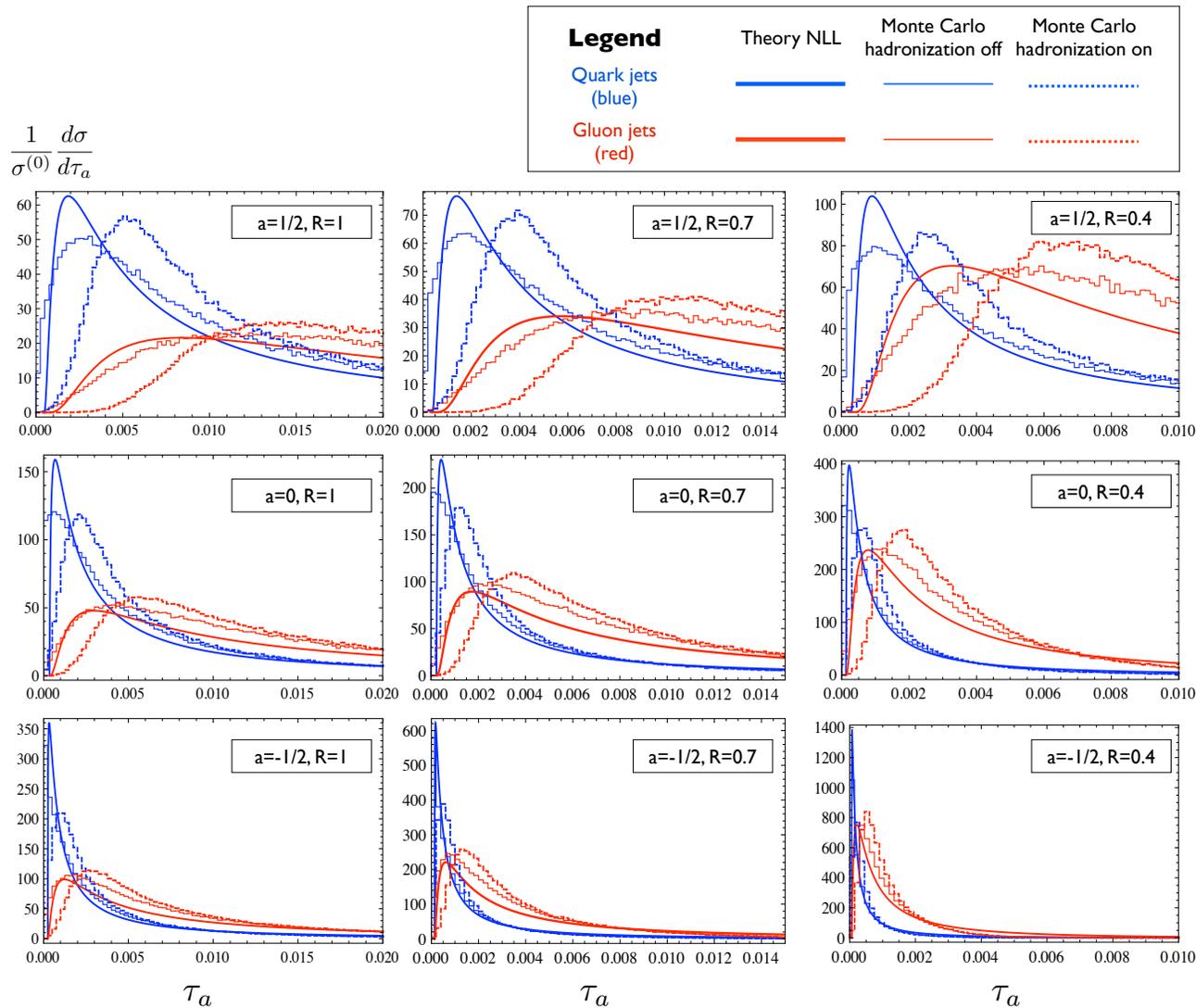
- Jet shapes in e^+e^- annihilation to 3 Jets:



- Measure one jet's shape, leave others unmeasured
- New factorization theorem to account for algorithm dependence (2010)

Resummed Predictions of Jet Shape Distributions vs. Monte Carlo

Ellis, Hornig, CL, Vermilion, Walsh



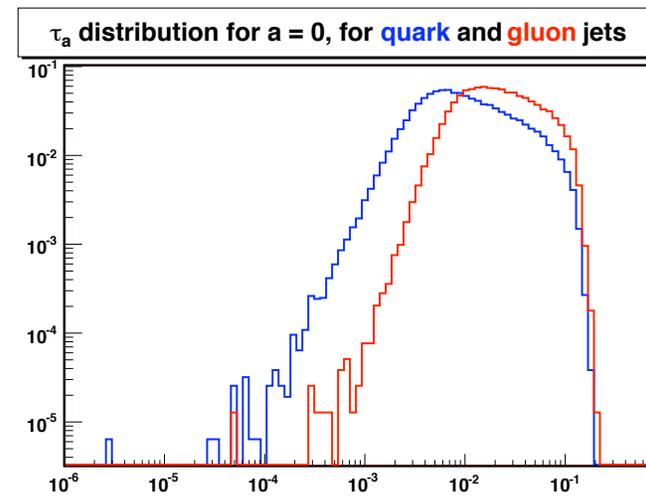
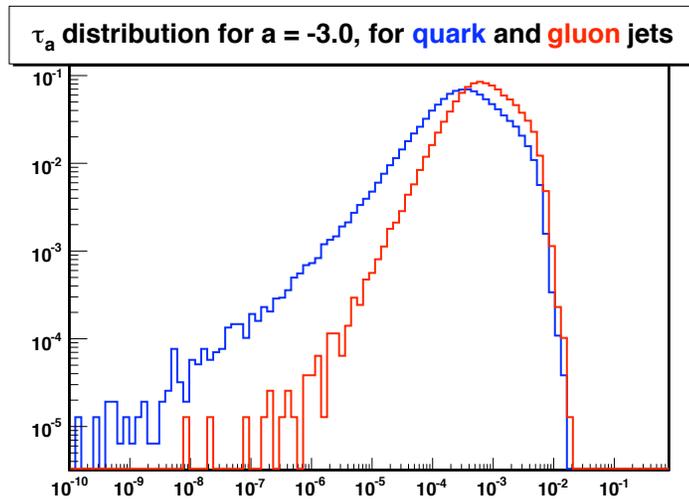
Q = 500 GeV
E_J = 150 GeV



Using angularities to distinguish quark and gluon jets

PRELIMINARY

Ellis, Hornig, CL, Vermilion, Walsh

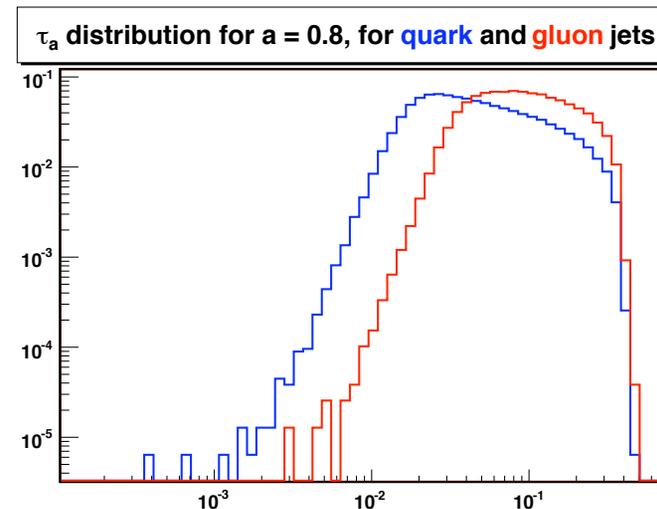


Studied quark v. gluon jets 3 well-separated jets in e^+e^-

cuts exist keeping $\sim 2\%$ of gluon jets and $\sim 20\%$ of quark jets (**10:1** enhancement)

or $\sim 15\%$ of gluons and 8% quarks (**2:1** enhancement).

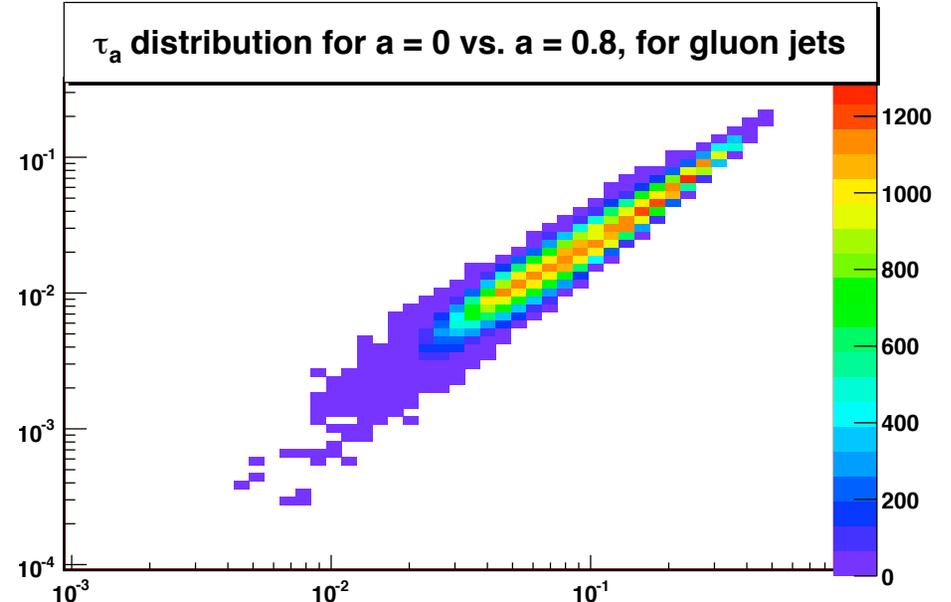
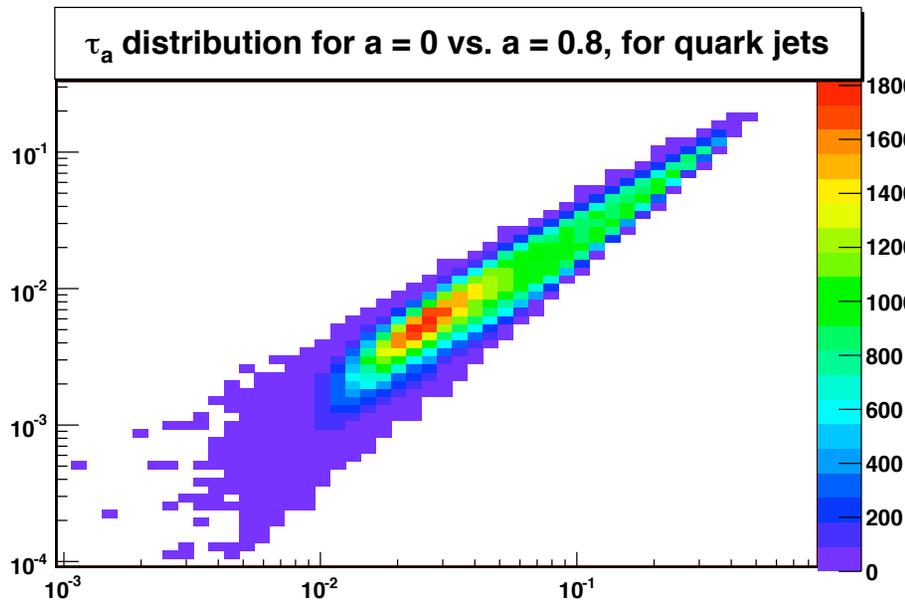
Expect greater discriminating power in correlated distributions for multiple values of a



Using angularities to distinguish quark and gluon jets

PRELIMINARY

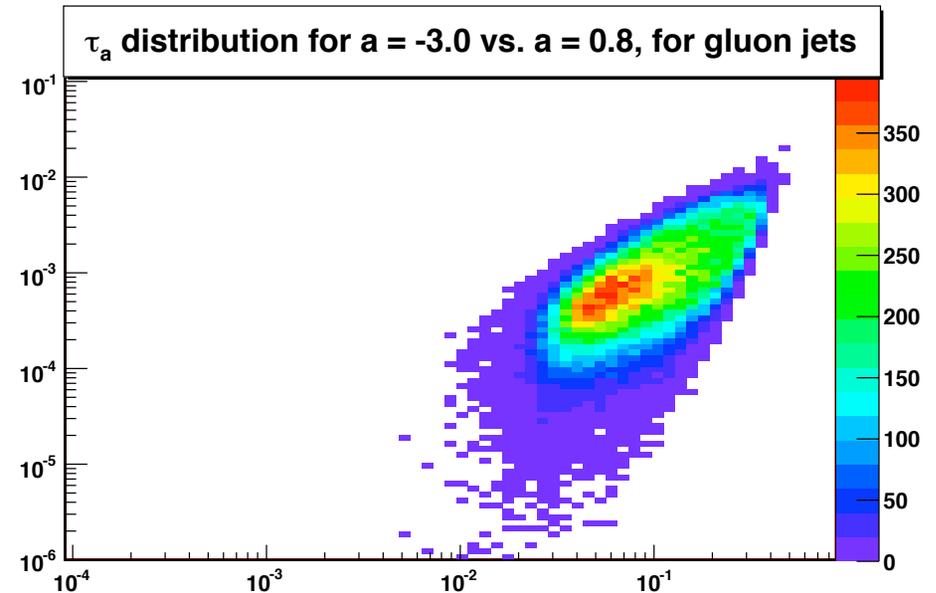
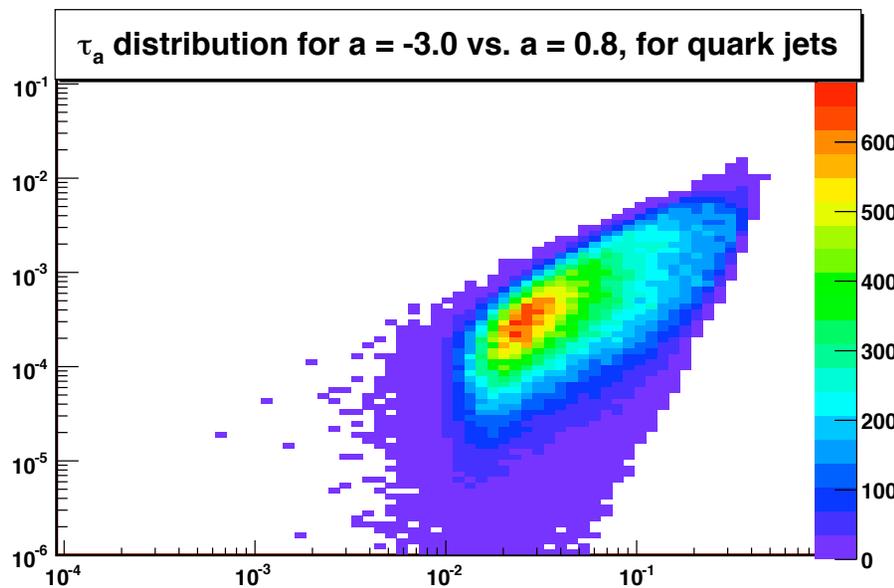
Ellis, Hornig, CL, Vermilion, Walsh



Using angularities to distinguish quark and gluon jets

PRELIMINARY

Ellis, Hornig, CL, Vermilion, Walsh



May find more elaborate cuts in these planes with greater distinguishing power than 1-D plots

Suggests use of multivariate analysis of many values of a to maximize discriminatory power.
(**really preliminary:** factor of **3-4** enhancement over single variable discriminatory power)

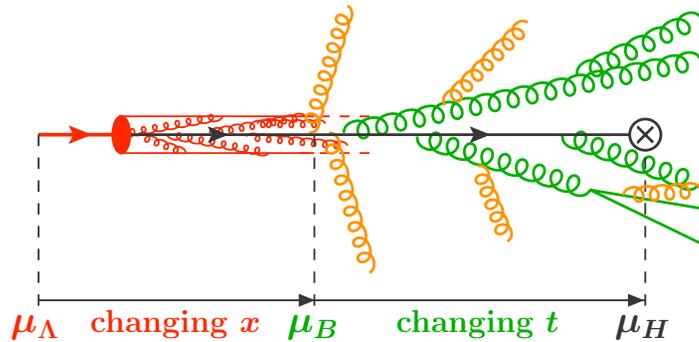
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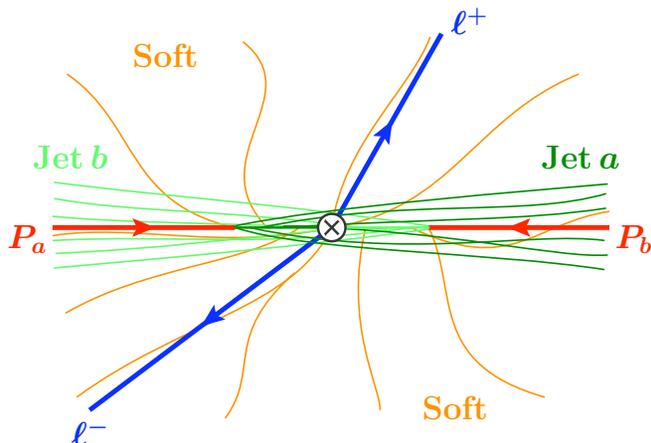


Progress in Describing Hadron Collisions

“Beam Functions”: PDFs + ISR

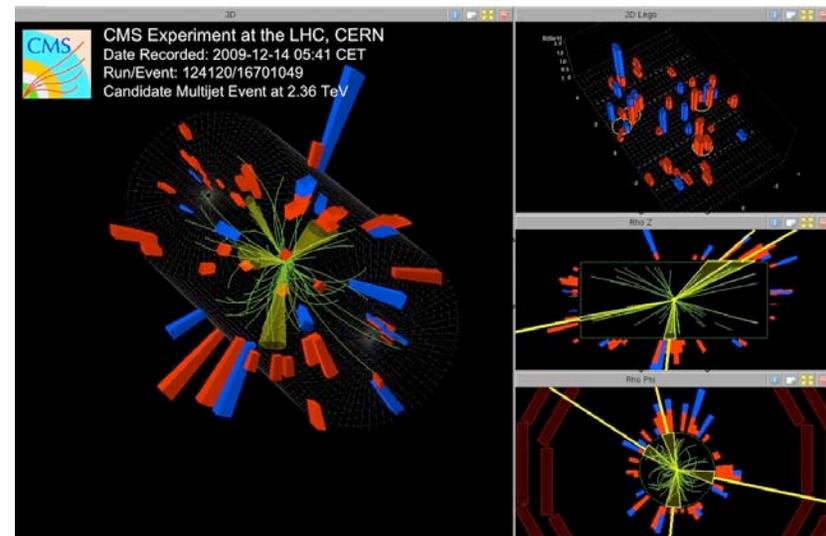


“Beam Thrust” in Drell-Yan



“N-jettiness” hadronic event shape

Stewart, Tackmann, Waalewijn



Top quark pair production

Ahrens, Ferroglia, Neubert, Pecjak, Yang

Higgs production

Chiu, Fuhrer, Kelley, Manohar; Ahrens, Becher, Neubert, Yang; Mantry, Petriello

“Dynamical” threshold enhancements from PDF steepness

Becher, Neubert, Xu; Bauer, Dunn, Hornig

Jets in Dense QCD Media

Idillbi, Majumder; D’Eramo, Liu, Rajagopal

New Tools for Probing Jets at LHC



- **A unified framework** in EFT for **factorization**, **resummation**, and **universal nonperturbative corrections** to jet shape distributions and other jet cross-sections
- **Recent new tools and applications of SCET to jets**
 - first consistent factorization of jet shape distributions with jet algorithms
 - first NLL resummation of jet angularity jet shapes
 - most precise extraction of strong coupling with consistent treatment of NP effects
 - hadronic event shapes, top & Higgs cross-sections at LHC
- **Future directions and challenges**
 - Distinguishing quark, gluon, and other jets using jet shapes and substructure
 - Resumming logarithms induced by phase space cuts from jet algorithms
 - Devising and predicting more observables suitable at LHC
- **2010+: EFT and Jet Shapes will help distinguish SM and New Physics at the LHC**